

# ELECTROMAGNETIC WAVES

## CHAPTER 33

Huygens - Light is a wave.

---

Einstein - Light is a particle.

---

Maxwell - Light is a  
Electromagnetic wave.

\*frustrated  
students



- Electromagnetic Oscillations & Alternating Current
- Maxwell's Equations & Magnetism of Matter
- **Electromagnetic Waves**
- Images
- Interference
- Diffraction

## Videos links:

[How Did We Figure Out What Light Is?](#) (history of light)

[Light Is Waves: Crash Course Physics #39](#)

[Spectra Interference: Crash Course Physics #40](#)

[The origin of Electromagnetic waves, and why they behave as they do](#)

[Understanding Electromagnetic Radiation!](#)

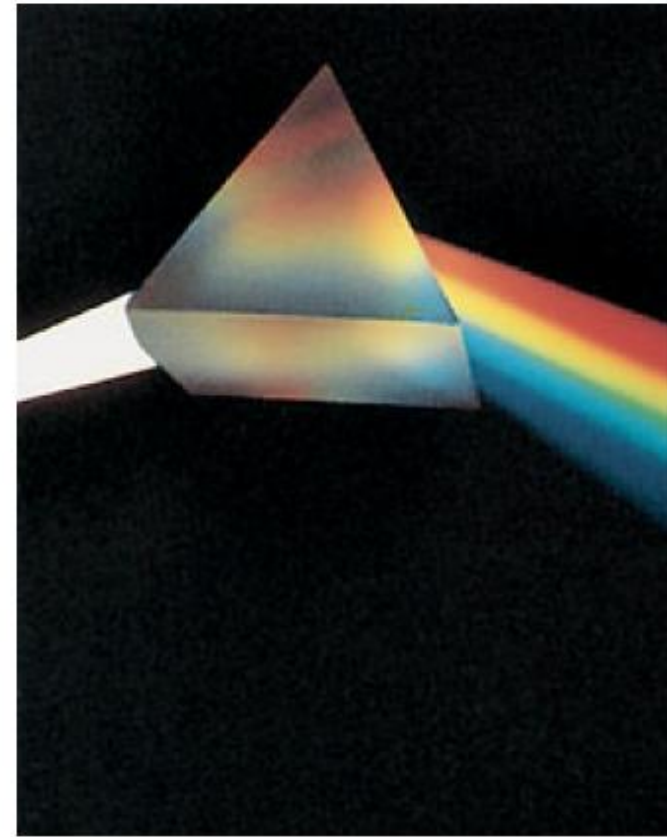
[But why would light "slow down"?](#)

# ELECTROMAGNETIC WAVES

Textbook: Chapter 33

- EM SPECTRUM AND TRAVELING EM WAVES
- ENERGY TRANSPORT AND POYNTING VECTOR
- RADIATION PRESSURE
- POLARIZATION
- REFLECTION AND REFRACTION
- TOTAL INTERNAL REFLECTION

Note: EM: Electromagnetic REFLECTION



Courtesy Bausch & Lomb

# EM SPECTRUM AND TRAVELING EM WAVES

James Clark Maxwell

→ **Light = Traveling EM wave**

**EM wave:**

**Electric Field  $\vec{E}$**

**+**

**Magnetic Field  $\vec{B}$**

Optics is Electromagnetism

Examples of EM waves

- Visible light
- Infrared light
- UV light
- $\mu$  waves
- Radio waves
- ...

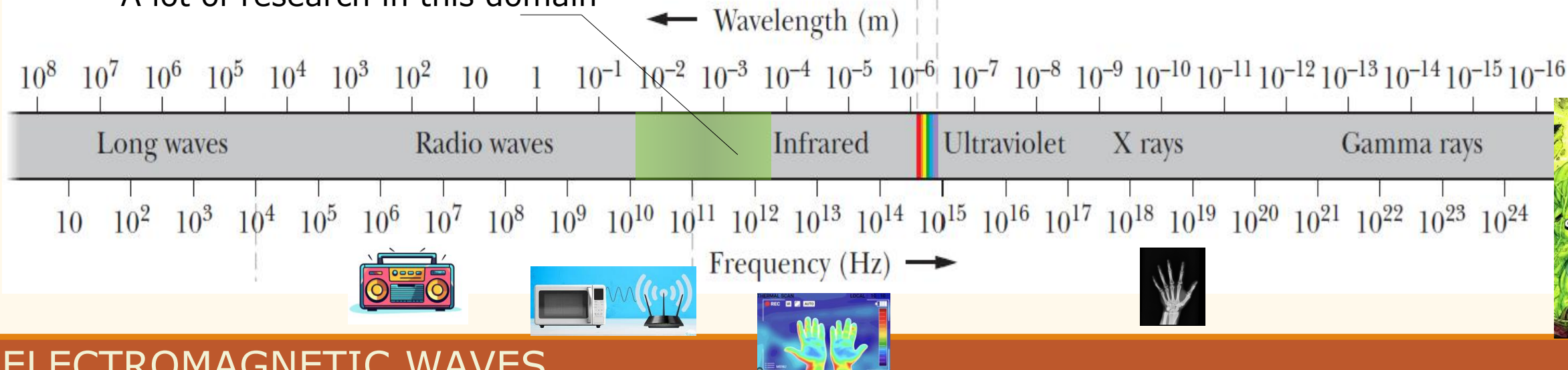
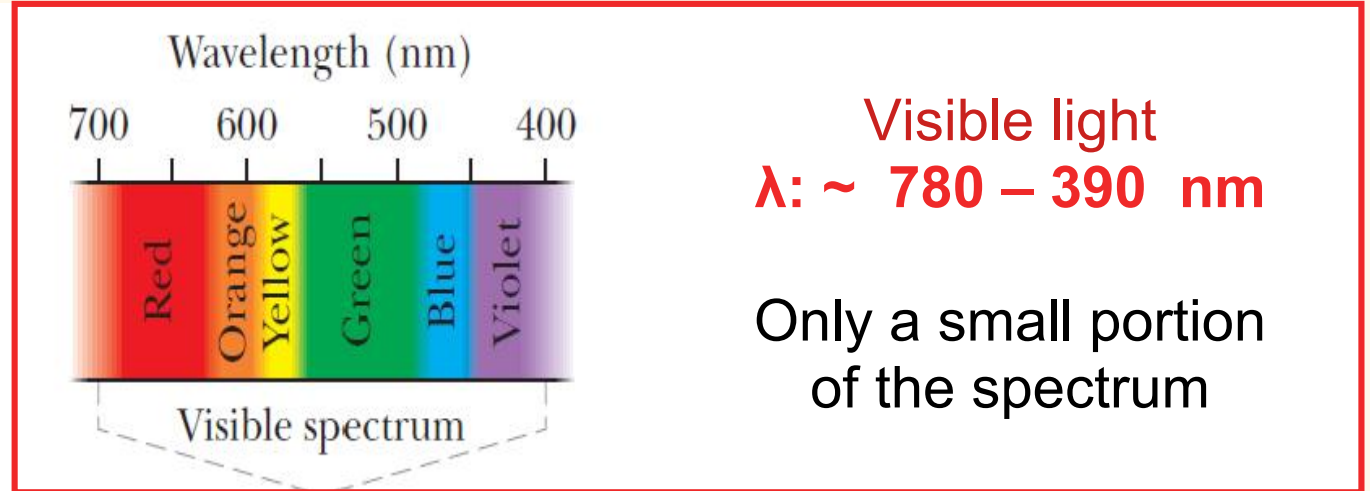
**EM can be sorted into domains**  
regarding their frequency

# EM SPECTRUM AND TRAVELING EM WAVES

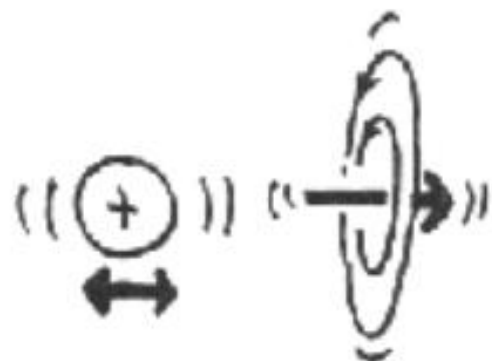
All EM waves travel in the **free space**  
(*vacuum*) at speed **c**

$$c = \frac{1}{\sqrt{\epsilon_0 \mu_0}} = 3 \cdot 10^8 \text{ m/s}$$

THz domain:  $\lambda$ :  $\sim 3 \text{ mm} - 30 \mu\text{m}$   
A lot of research in this domain



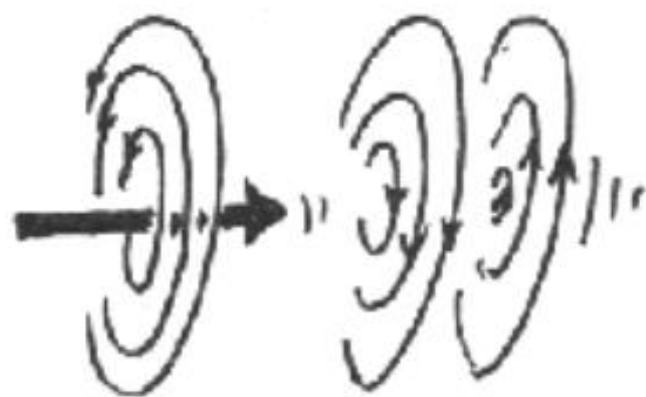
IMAGINE A SINGLE ELECTRIC CHARGE BEING **VIBRATED**:



IN THE SPACE NEAR THE VIBRATING CHARGE, THE CHARGE'S ELECTRIC FIELD IS CHANGING, SO IT INDUCES A MAGNETIC FIELD CURLING AROUND IT.

BUT THE MAGNETIC FIELD IS ALSO CHANGING — SO IT INDUCES MORE ELECTRIC FIELD, WHICH INDUCES MORE MAGNETIC FIELD...

**ETC.!**



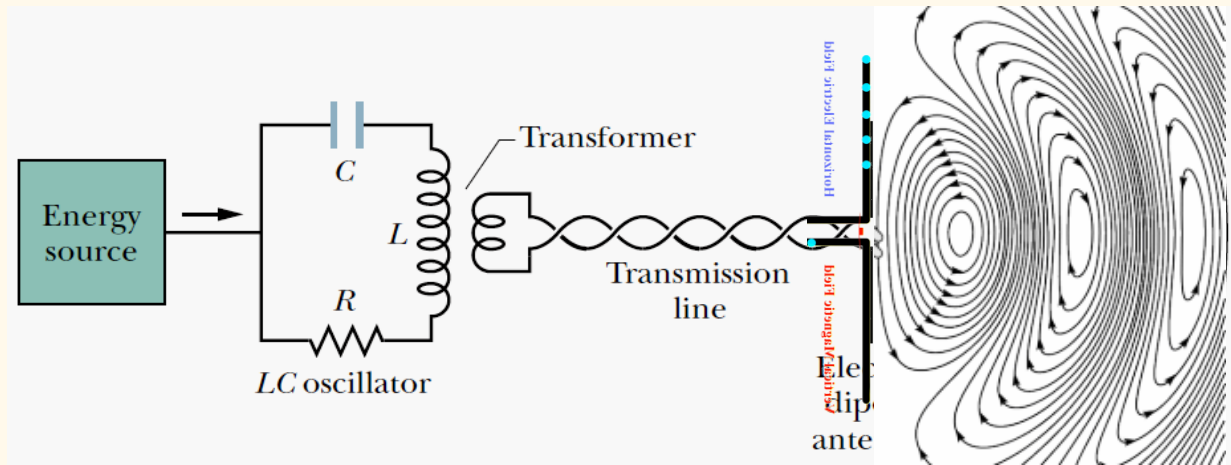
# EM SPECTRUM AND TRAVELING EM WAVES

Generation of EM

→ oscillating electric charges

Example for  $\lambda > 1\text{m}$  :  
LC oscillator coupled to a  
dipole antenna

$$\omega = \frac{1}{\sqrt{LC}}$$



Current  
oscillates at  $\omega$

Dipole  
oscillates at  $\omega$

Electric field and Magnetic field  
oscillate at  $\omega$

**Change over space  $\neq$  instantaneous  
EM wave of angular freq.  $\omega$   
traveling at speed  $c$**

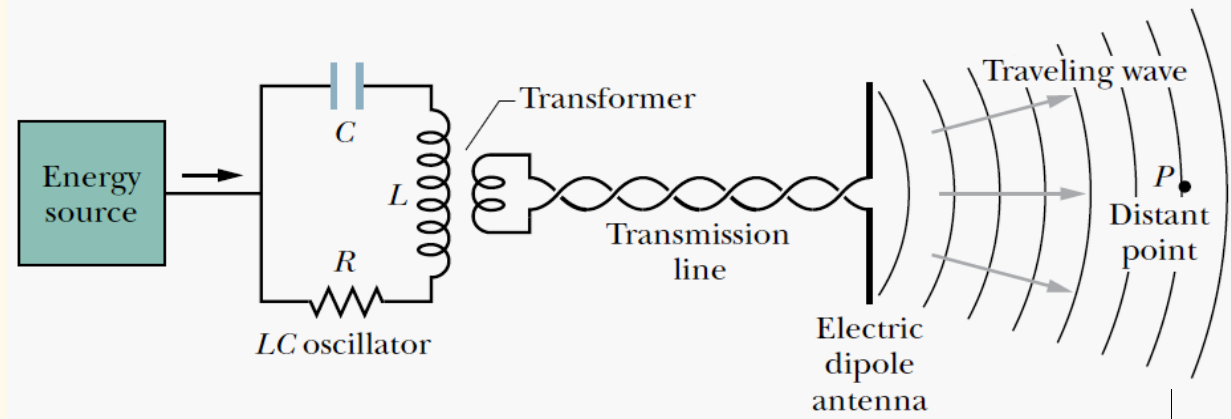


# EM SPECTRUM AND TRAVELING EM WAVES

## Structure of the EM wave

→ **Far from the antenna** so the *curvature of the wave* can be neglected: **Plane wave**

Out of the plane wave approximation physical description of EM waves is more complex



**Current**  
oscillates at  $\omega$

**Dipole**  
oscillates at  $\omega$

**Electric field** and **Magnetic field**  
oscillate at  $\omega$

**Change over space  $\neq$  instantaneous**  
**EM wave of angular freq.  $\omega$**   
**traveling at speed  $c$**

# EM SPECTRUM AND TRAVELING EM WAVES

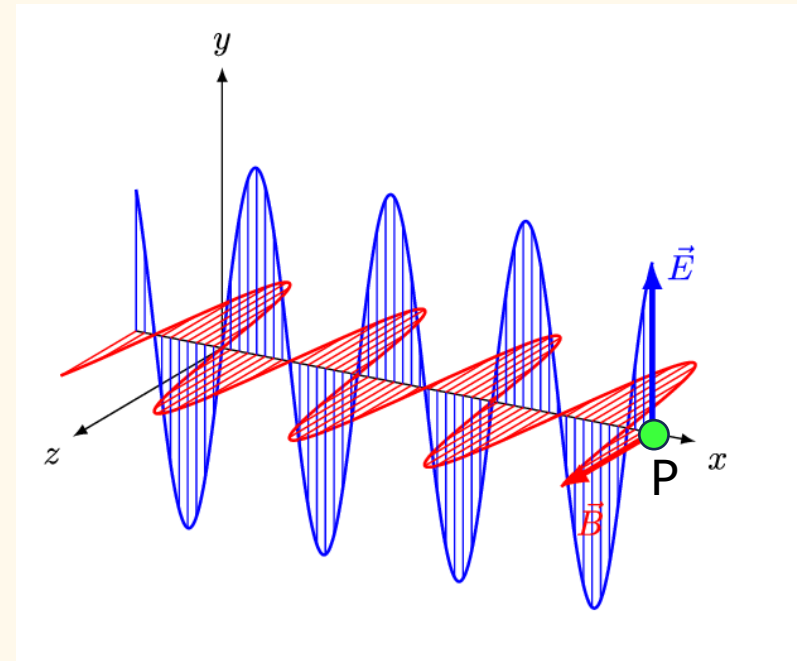
## Structure of the EM traveling plane wave

$\vec{E}$  and  $\vec{B}$  are **perpendicular to the direction of propagation** → **transverse wave**

$\vec{E}$  is **perpendicular** to  $\vec{B}$

$\vec{E} \times \vec{B}$  gives the **direction of propagation**

$\vec{E}$  and  $\vec{B}$  vary sinusoidally at the **same  $\omega$  in phase**



P fixed point in space

# EM SPECTRUM AND TRAVELING EM WAVES

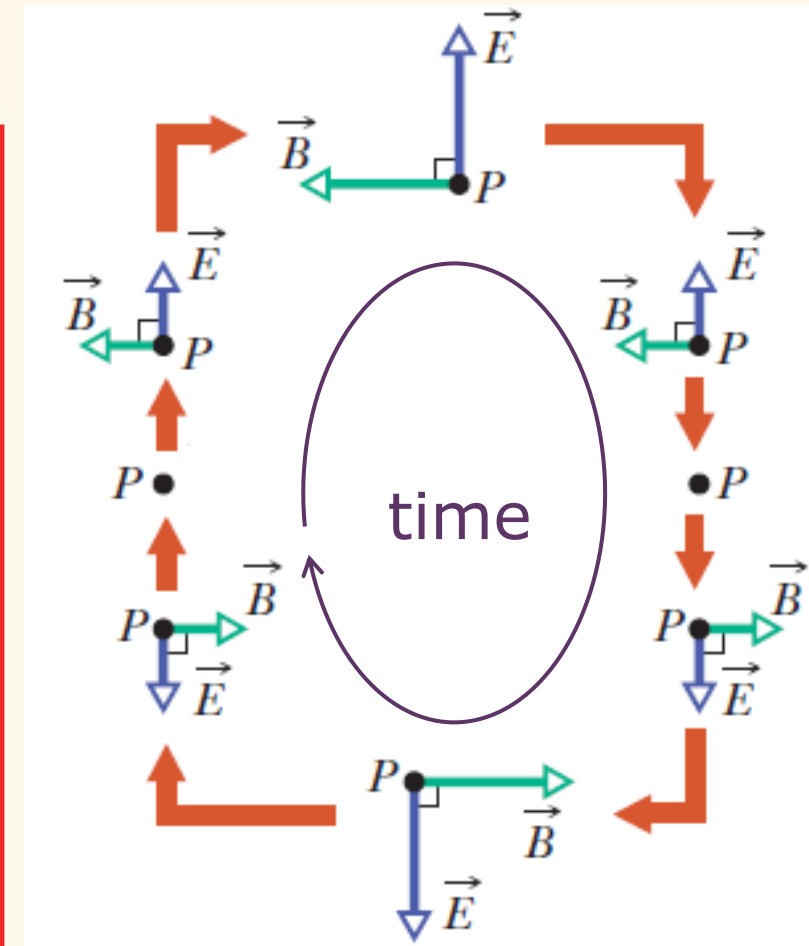
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# EM SPECTRUM AND TRAVELING EM WAVES

## Structure of the EM traveling plane wave

$\vec{E}$  and  $\vec{B}$  vary sinusoidally at the **same  $\omega$  in phase**

Assuming **forward propagation** along the **x** axis

$$E = E_m \sin(kx - \omega t)$$

$$B = B_m \sin(kx - \omega t)$$

$E_m, B_m$ : amplitudes of E and B

$k$ : angular wave number

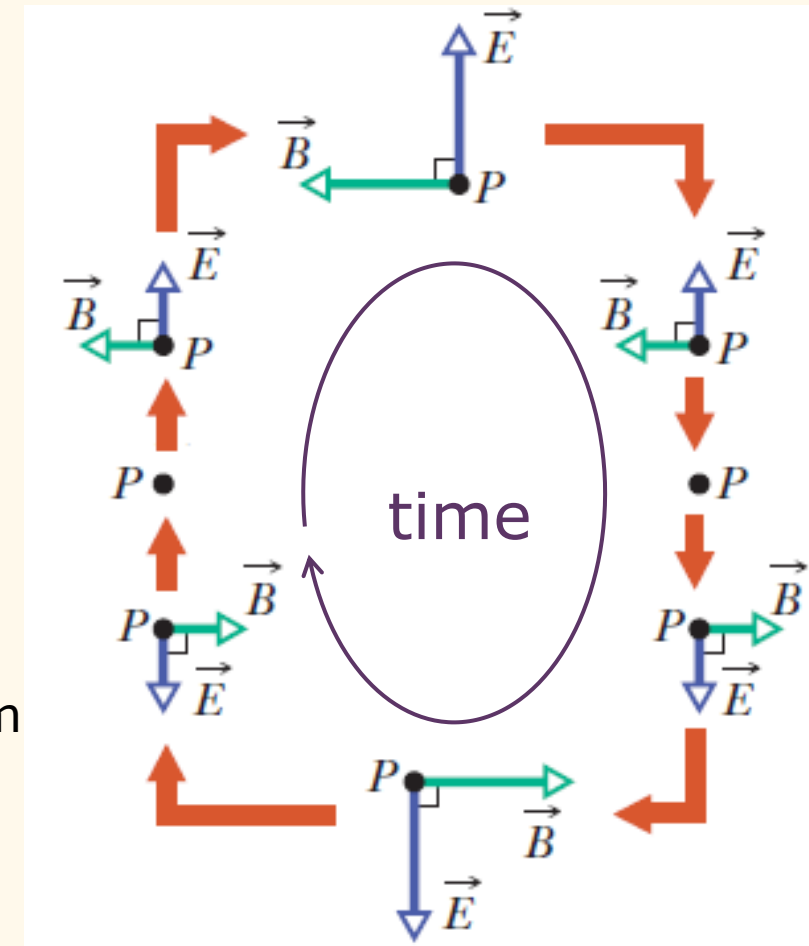
$\omega$ : angular frequency

$c = \omega / k$  speed of light in vacuum

Valid for linear or circular polarization

$$\frac{E_m}{B_m} = \frac{E}{B} = c$$

Amplitude & magnitude ratios

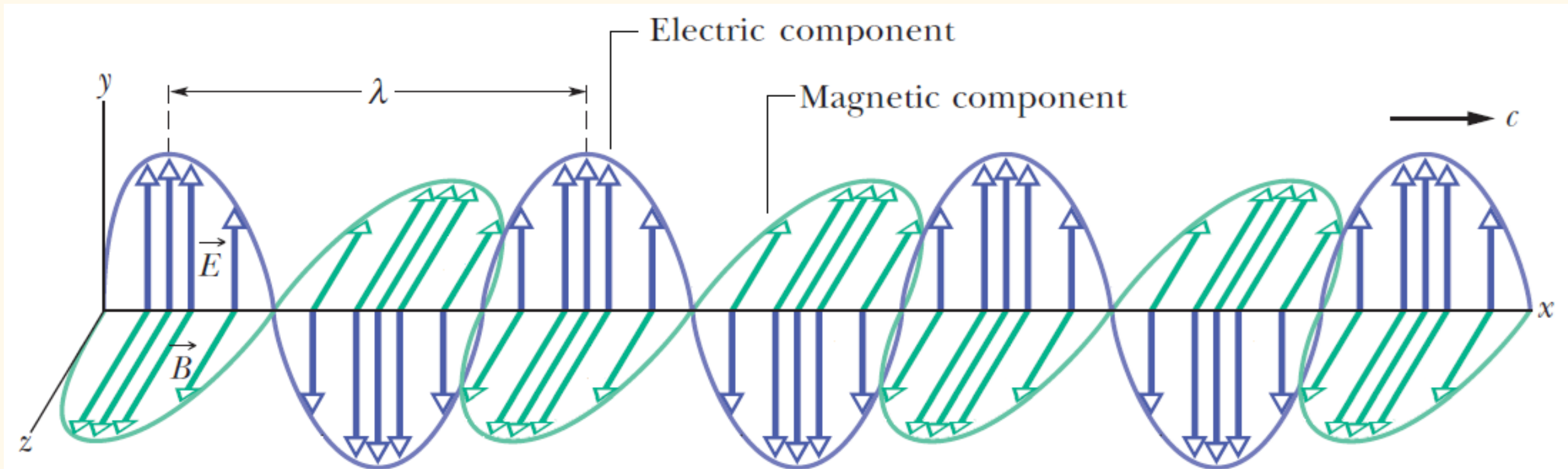


P fixed point in space

# EM SPECTRUM AND TRAVELING EM WAVES

Representation of the wave:

Representations of the **fields** through space at a **fixed instant**



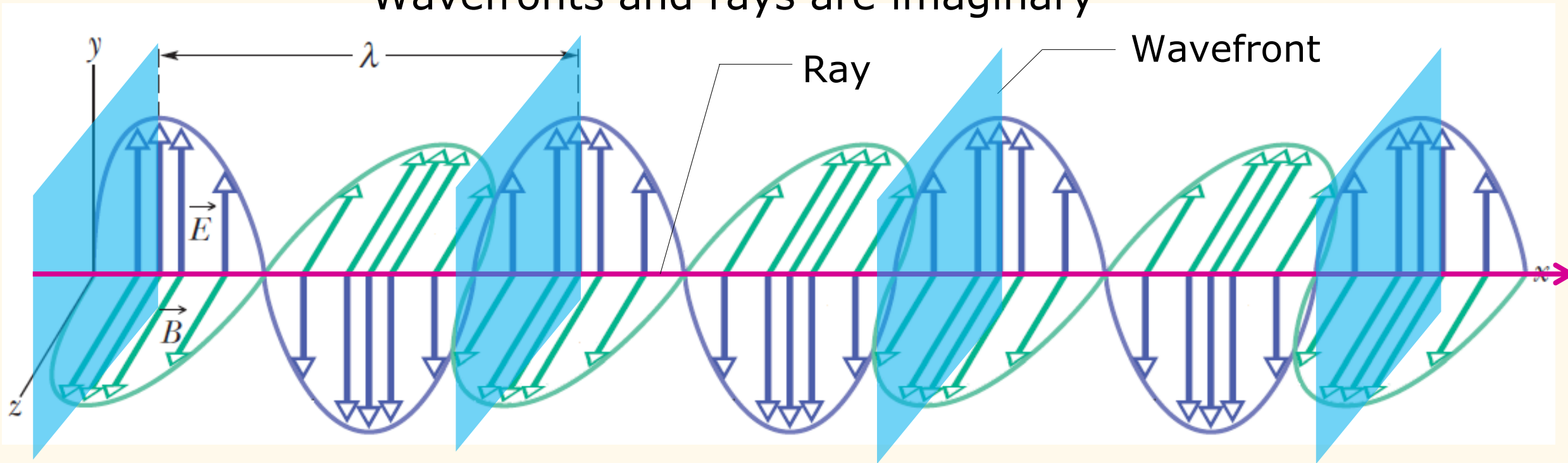
# EM SPECTRUM AND TRAVELING EM WAVES

Representation of the wave:

**Wavefronts:** Surfaces with  $E$  at **constant phase** separated by  $\lambda$

**Rays:** Line in the **direction of propagation** perpendicular to the wavefront

Wavefronts and rays are imaginary



# EM SPECTRUM AND TRAVELING EM WAVES

## Specificity of EM waves

- Unlike other waves (e.g. sound) EM propagate **without medium**
- Special Relativity Theory:

The speed of light  $c$  is **the same no matter the motion** of the source or the observer

## Propagation of EM waves

- Varying B field induces a varying E field (M-F)
- Varying E field induces a varying B field (M-A)
- And so on ...

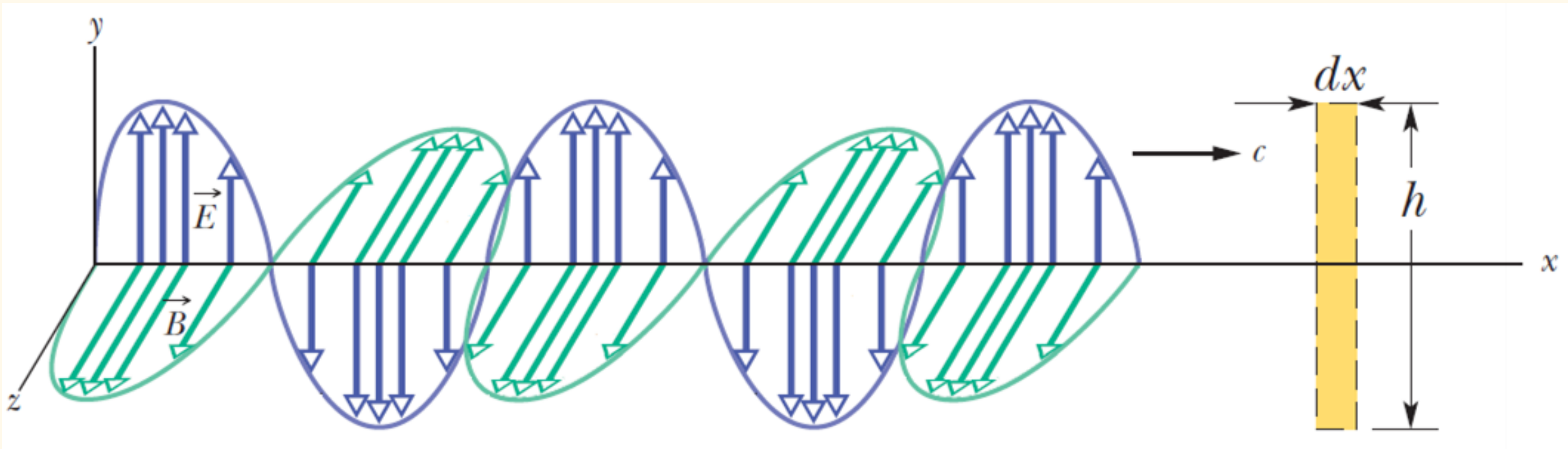
The two fields **continuously create each other** via induction

→ Next: Exploring this statement

# EM SPECTRUM AND TRAVELING EM WAVES

## Quantitative analysis

We consider a small rectangle ( surface =  $h \, dx$ ) perpendicular to Oz





# EM SPECTRUM AND TRAVELING EM WAVES

## Quantitative analysis

During  $dt$   $\vec{B}$  has varied  $\rightarrow \phi_B$  in the rectangle varies  $\rightarrow$  induced  $\vec{E}$  field

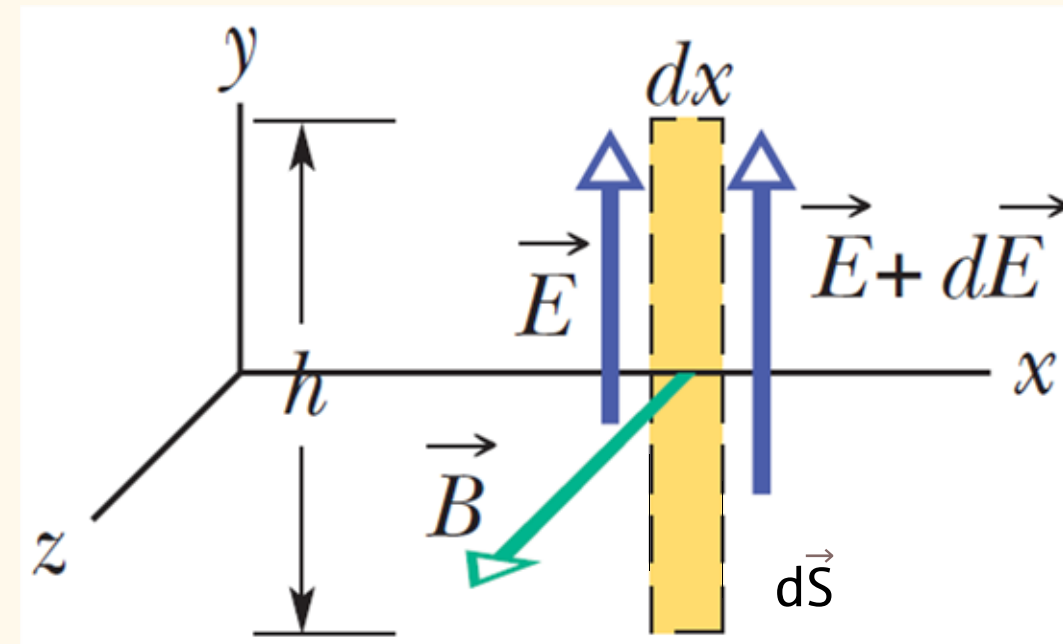
On one edge electric field is  $\vec{E}$  and  $\vec{E} + d\vec{E}$  on the other side

$$\oint \vec{E} \cdot d\vec{S} = -\frac{d\phi_B}{dt} \quad (M - F)$$

$$0 + (\cancel{E} + dE)\cancel{h} + 0 - \cancel{E}h = -\frac{d}{dt}(B \cancel{h} dx)$$

$$\frac{dE}{dx} = -\frac{dB}{dt} \quad \boxed{\frac{\partial E}{\partial x} = -\frac{\partial B}{\partial t}}$$

Note: 2 variables  $\rightarrow$  partial derivative



# EM SPECTRUM AND TRAVELING EM WAVES

## Quantitative analysis


$$\frac{\partial E}{\partial x} = -\frac{\partial B}{\partial t}$$

$$E = E_m \sin(kx - \omega t)$$

$$B = B_m \sin(kx - \omega t)$$

$$\frac{\partial E}{\partial x} = E_m k \cos(kx - \omega t)$$

$$\frac{\partial B}{\partial t} = -B_m \omega \cos(kx - \omega t)$$


$$E_m k \cos(kx - \omega t) = B_m \omega \cos(kx - \omega t)$$

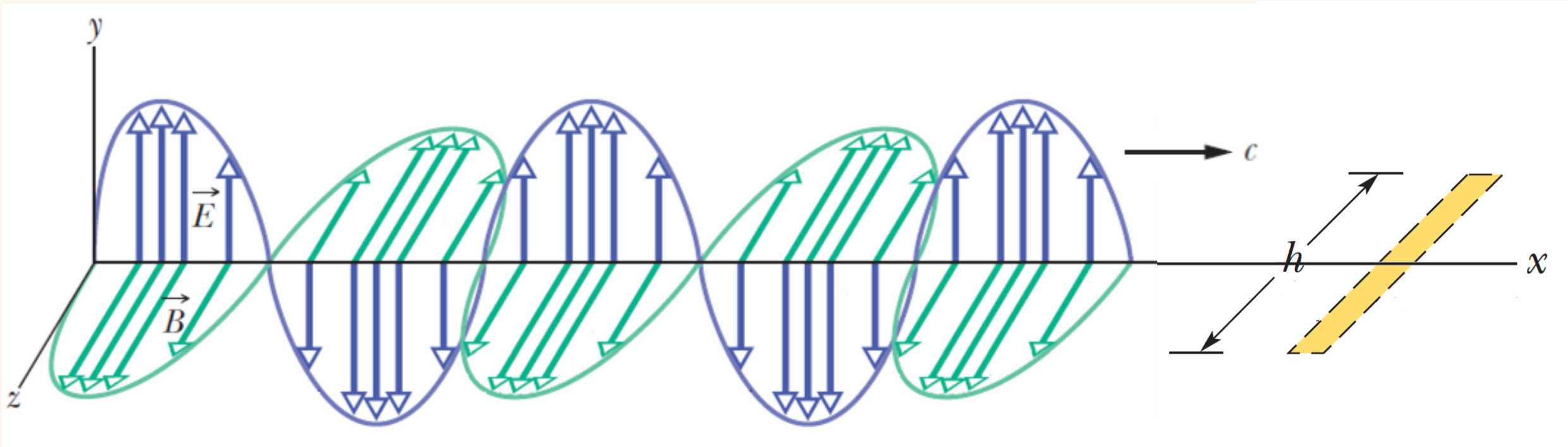
That demonstrates

$$\frac{E_m}{B_m} = \frac{\omega}{k} = c$$

# EM SPECTRUM AND TRAVELING EM WAVES

## Quantitative analysis

We consider a small rectangle ( surface =  $h \, dx$ ) perpendicular to Oy



# EM SPECTRUM AND TRAVELING EM WAVES

## Quantitative analysis

During  $dt$   $\vec{E}$  has varied  $\rightarrow \phi_E$  in the rectangle varies  $\rightarrow$  induced B field

On one edge electric field is  $\vec{B}$  and  $\vec{B} + d\vec{B}$  on the other side

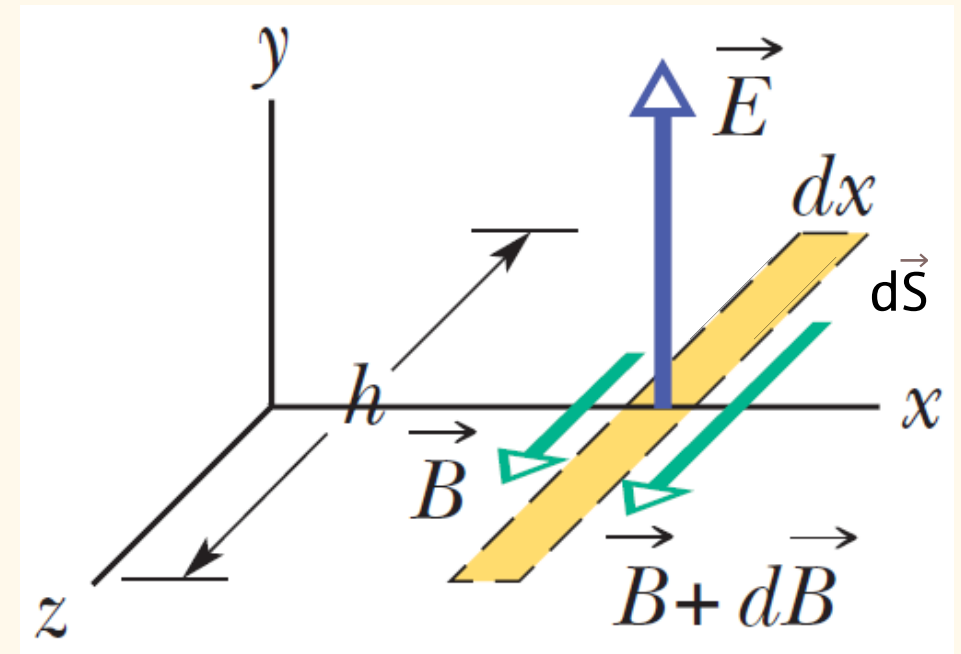
$$\oint \vec{B} \cdot d\vec{S} = \mu_0 \epsilon_0 \frac{d\phi_E}{dt} + \mu_0 i_{enc} \quad (\text{M - A})$$

Note: we assume free space propagation:  $i_{enc} = 0$

$$0 - (\cancel{B} + dB)\cancel{h} + 0 + \cancel{B}h = \mu_0 \epsilon_0 \frac{d}{dt} (E \cancel{h} dx)$$

$$-\frac{dB}{dx} = \mu_0 \epsilon_0 \frac{dE}{dt} \quad \boxed{-\frac{\partial B}{\partial x} = \mu_0 \epsilon_0 \frac{\partial E}{\partial t}}$$

Note: 2 variables  $\rightarrow$  partial derivative



# EM SPECTRUM AND TRAVELING EM WAVES

## Quantitative analysis


$$-\frac{\partial B}{\partial x} = \mu_0 \epsilon_0 \frac{\partial E}{\partial t}$$

$$E = E_m \sin(kx - \omega t)$$

$$B = B_m \sin(kx - \omega t)$$

$$\frac{\partial E}{\partial t} = -E_m \omega \cos(kx - \omega t)$$

$$\frac{\partial B}{\partial x} = B_m k \cos(kx - \omega t)$$


$$B_m k \cos(kx - \omega t) = \mu_0 \epsilon_0 E_m \omega \cos(kx - \omega t)$$

$$\frac{E_m \omega}{B_m k} = \frac{1}{\mu_0 \epsilon_0}$$

But also  $\frac{E_m}{B_m} = \frac{\omega}{k} = c$  —————  $c^2 = \frac{1}{\mu_0 \epsilon_0}$  —————

$$c = \frac{1}{\sqrt{\mu_0 \epsilon_0}}$$

# ENERGY TRANSPORT AND POYNTING VECTOR

$$\vec{S} = \frac{\vec{E} \times \vec{B}}{\mu_0}$$

Since  $\vec{E}$  and  $\vec{B}$  are perpendicular:

- $\vec{S}$  is along the direction of propagation

- $S = \frac{EB}{\mu_0} = \frac{1}{c\mu_0} E^2$

## EM waves carry energy

Characterized by the  
**Poynting vector**  $\vec{S}$

$S$  relates to the instantaneous  
energy transfer rate per area

$$S = \left( \frac{\text{energy/time}}{\text{surface}} \right)_{inst} (W / m^2)$$

Note: Poynting vector  $\neq$  spin

# ENERGY TRANSPORT AND POYNTING VECTOR

$$I = S_{avg} = \left( \frac{1}{c\mu_0} E^2 \right)_{avg} = \frac{1}{c\mu_0} (E^2)_{avg}$$

$$E = E_m \sin(kx - \omega t)$$

$$E^2 = E_m^2 \sin^2(kx - \omega t)$$

$$E^2_{avg} = E_m^2 (\sin^2(kx - \omega t))_{avg}$$

$$E^2_{avg} = E_m^2 / 2$$

$$I = \frac{E_m^2}{2c\mu_0}$$

## EM waves carry energy

Characterized by the  
**Poynting vector**  $\vec{S}$

Power transfer characterized by  
the **intensity I**

$$I = S_{avg} = \left( \frac{\text{energy / time}}{\text{surface}} \right)_{avg}$$
$$= \left( \frac{\text{power}}{\text{surface}} \right)_{avg}$$

Note: we usually find the notation :  $\langle S \rangle = (S)_{avg}$

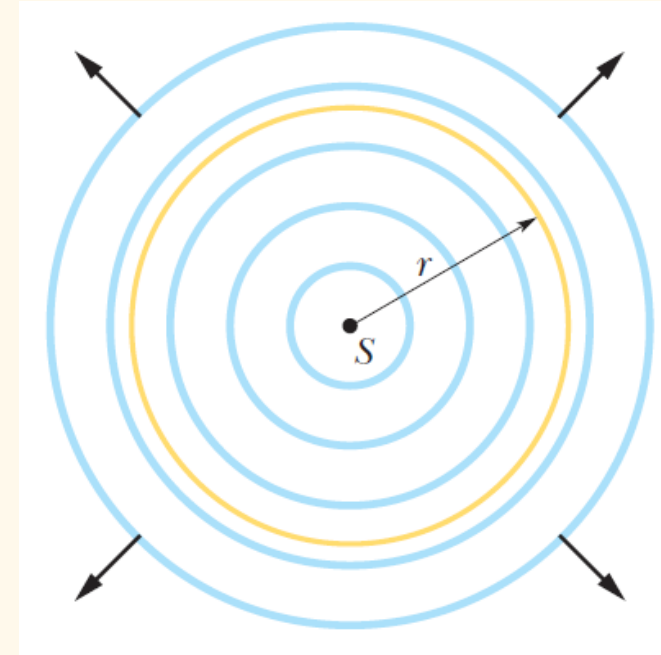
# EM SPECTRUM AND TRAVELING EM WAVES

## Variation of intensity with distance

We assume the EM comes from a point source of power  $P_s$   
→ spherical wavefronts close to the source

$$I = \frac{\text{power}}{\text{area}} = \frac{P_s}{4\pi r^2}$$

Intensity decreases with the squared distance from the source





# RADIATION PRESSURE

**Light** does not have mass but **carries momentum**

according to quantum mechanics  $p = \frac{U}{c}$  *Energy carry by the beam of light*  
*speed of light*



**EM waves carry energy**

**Energy & momentum**  
transfer to an irradiated object

# RADIATION PRESSURE

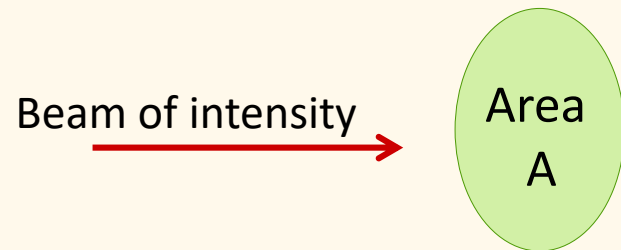
**Light** does not have mass but **carries momentum**

according to quantum mechanics  $p = \frac{U}{c}$  Energy carry by the beam of light  
speed of light

$$\Delta p = \frac{\Delta U}{c} \quad (\text{total absorption of light})$$

$$\Delta p = 2 \frac{\Delta U}{c} \quad (\text{total reflection of light})$$

$\Delta \vec{p}$  in direction of propagation of the incoming light



$$\Delta U = IA\Delta t$$

$$F = \frac{\Delta p}{\Delta t}$$

## EM waves carry energy

### Energy & momentum

transfer to an irradiated object

→ Force on the object

$$F = \frac{\Delta p}{\Delta t} = \frac{IA}{c} \quad (\text{total absorption})$$

$$F = 2 \frac{IA}{c} \quad (\text{total reflection})$$

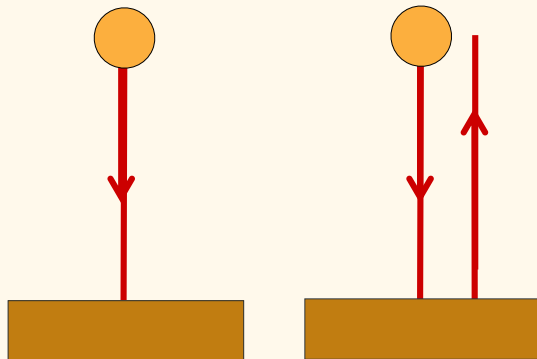
# RADIATION PRESSURE

$$\Delta p = \frac{\Delta U}{c} \quad F = \frac{IA}{c} \quad (\text{total absorption of light})$$

$$\Delta p = 2\frac{\Delta U}{c} \quad F = 2\frac{IA}{c} \quad (\text{total reflection of light})$$

Factor 2 between total absorption and total reflection

→ More momentum is transferred during an elastic collision rather than an inelastic collision



Note: if not total absorption or total reflection  
The factor is between 1 and 2

# RADIATION PRESSURE

$$I = \frac{\text{power}}{\text{area}} = \frac{\text{rate of doing work}}{\text{area}}$$

$W = \int F \cdot dl = F \cdot \Delta x$  in 1D and conservative force

$$= \frac{(F \cdot \Delta x) / \Delta t}{\text{area}}$$

$\Delta x / \Delta t = \text{speed} = c$   
in case of light

$$I = P_r c$$

$\frac{F}{\text{area}} = \text{pressure} = p_r$

$$p_r = \frac{I}{c}$$

also called "radiation pressure"

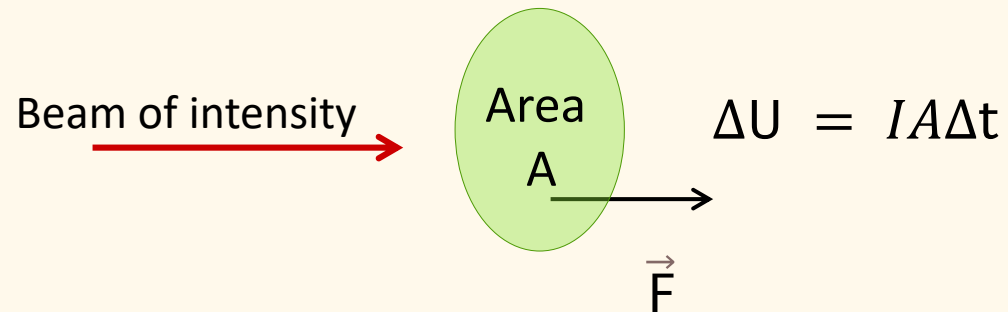
Note: Be careful not to confuse the symbol  $p_r$  for radiation pressure with the symbol  $p$  for momentum.

# RADIATION PRESSURE

$$P_r = \frac{I}{c} \quad (\text{total absorption of light})$$

$$P_r = 2 \frac{I}{c} \quad (\text{total reflection of light})$$

If  $I \gg$  and  $A \ll$   $\rightarrow$  (focused laser)  $P_r \gg$



**EM waves carry energy**

**Energy & momentum**  
transfer to an irradiated object

$\rightarrow$  Force on the object  
 $\rightarrow$  Pressure ( $P_r$ ) on the object

$$P_r = \frac{F}{A}$$

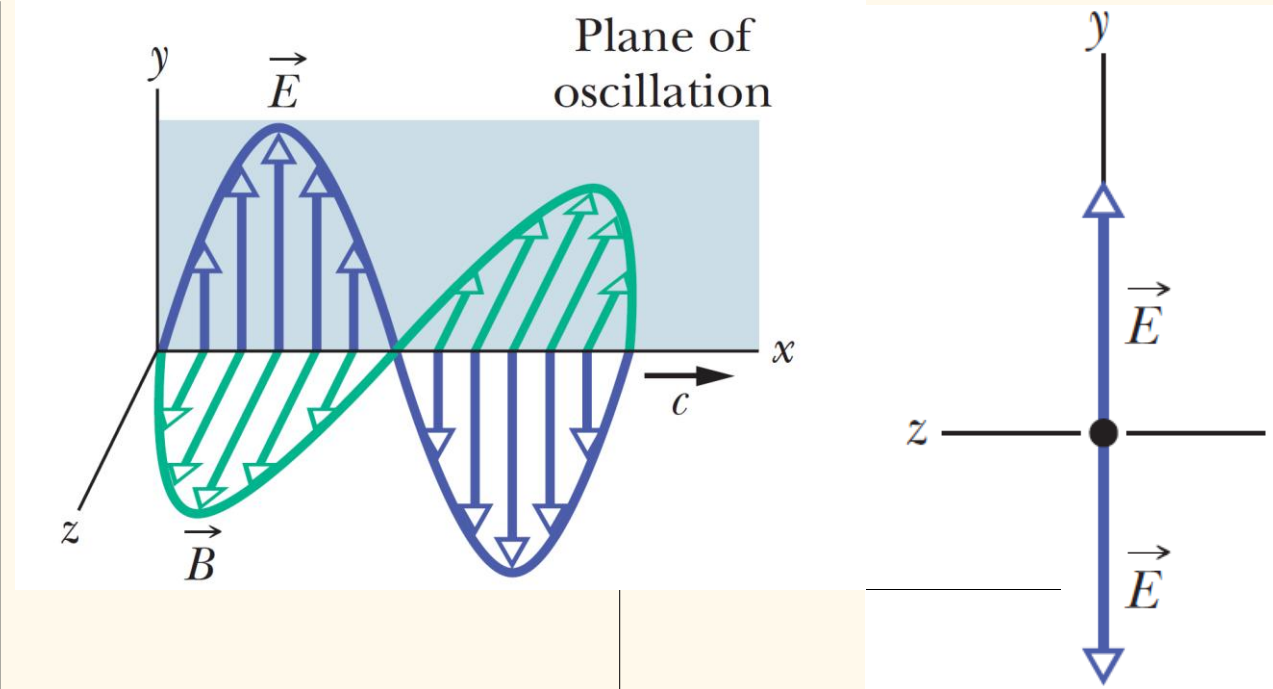
**Radiation Pressure**

# POLARIZATION

So far we represented the EM wave like this:

- $\vec{E}$  is always in the (Oxy) plane
- **Linear polarization** along Oy
- Vertical polarization
- (Oxy): **plane of oscillation**

**But light is not always linearly polarized**



Schematic representation of an EM wave polarized along Oy

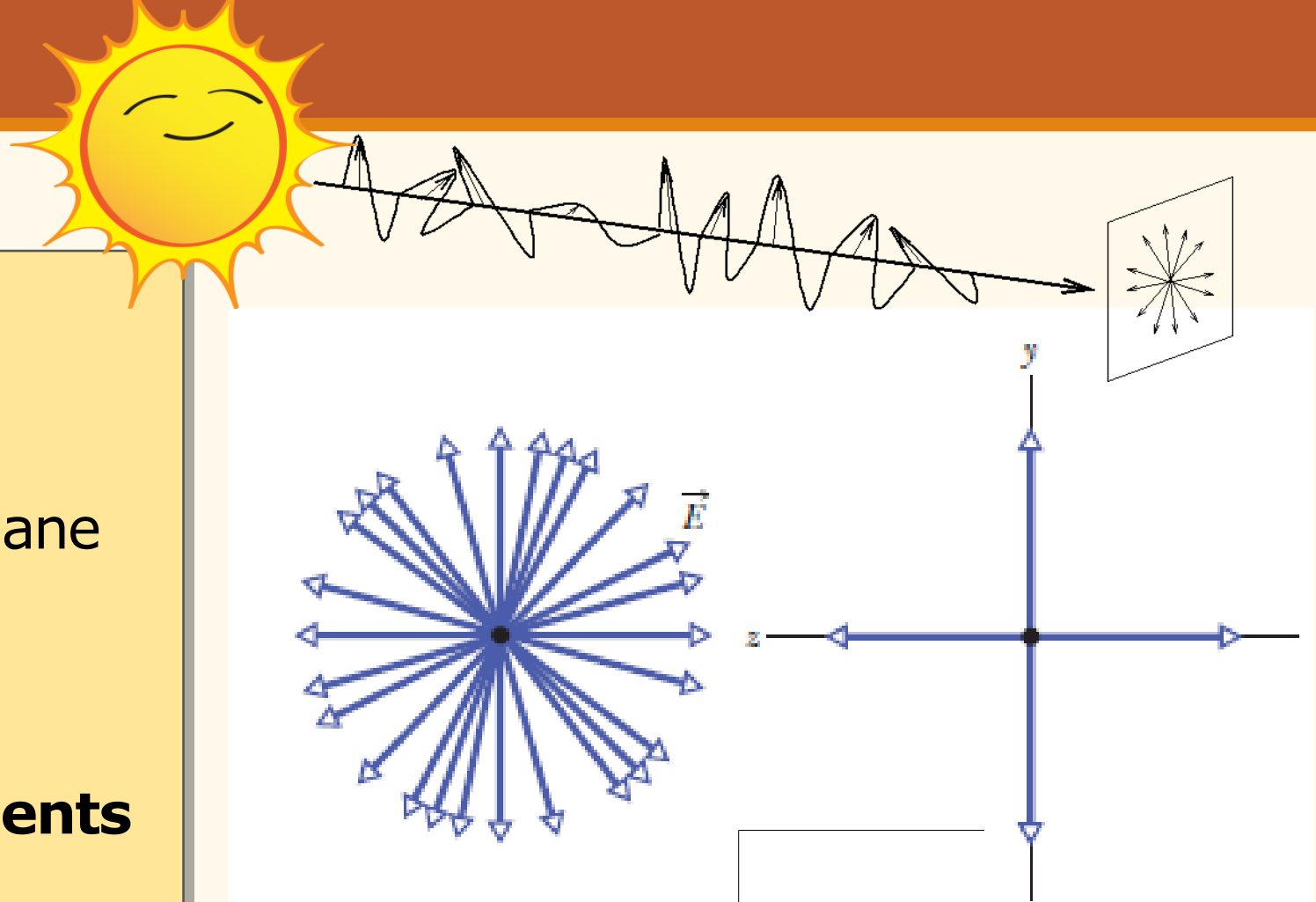
# POLARIZATION

**Unpolarized** light  
(e.g. sun, lightbulb, ... )

$\vec{E}$  **randomly oriented** in a plane  
perpendicular to the direction  
of propagation

Separation of  $\vec{E}$  in **2 components**  
→ along Oz & Oy:  $E_z$  and  $E_y$

**Both oscillate as linearly  
polarized electric fields**

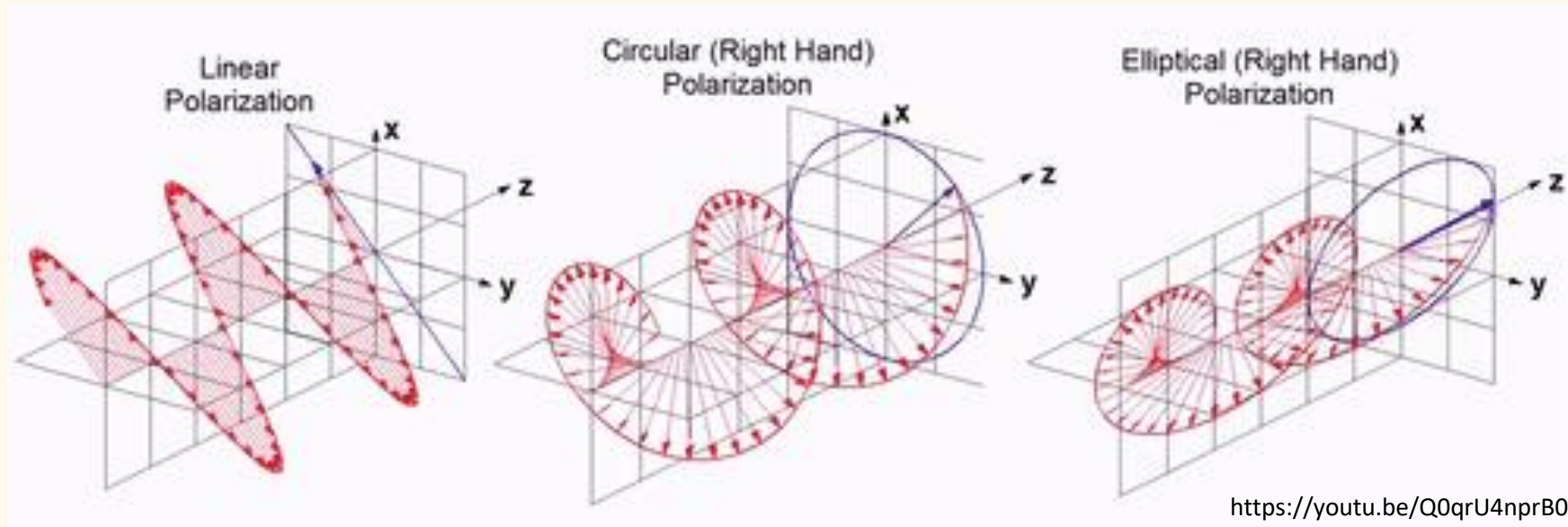
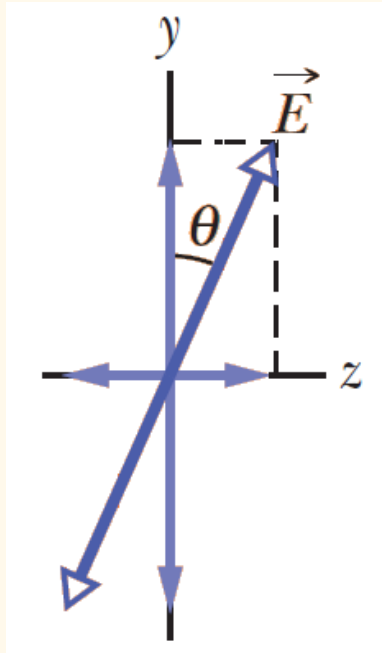


Schematic representation of  
an unpolarized EM wave

# POLARIZATION

## Note on polarization:

- It is possible to have a **partially polarized** light



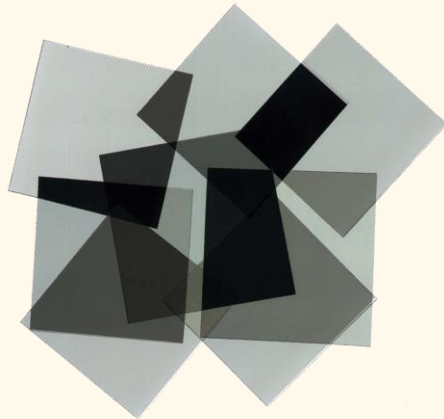
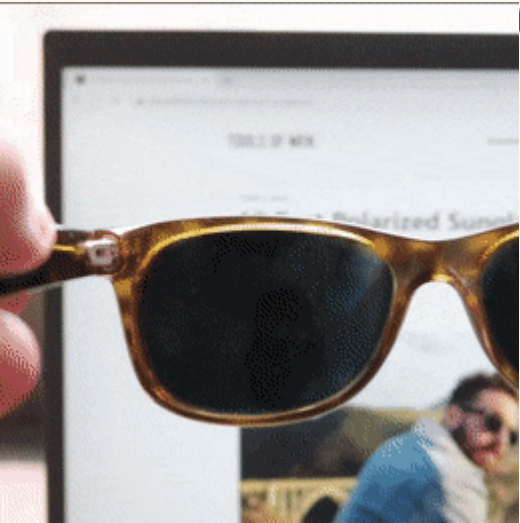
- If both components are in phase  
→ back to **linear polarization**  
oriented with angle  $\theta$  with respect to Oy

- If  $E_z$  and  $E_y$  have a  $\pm\pi/2$  phase shift and the same amplitude  
→ **circular polarization**

- If  $E_z$  and  $E_y$  have a  $\pm\pi/2$  phase shift  
→ **elliptic polarization**



# POLARIZATION

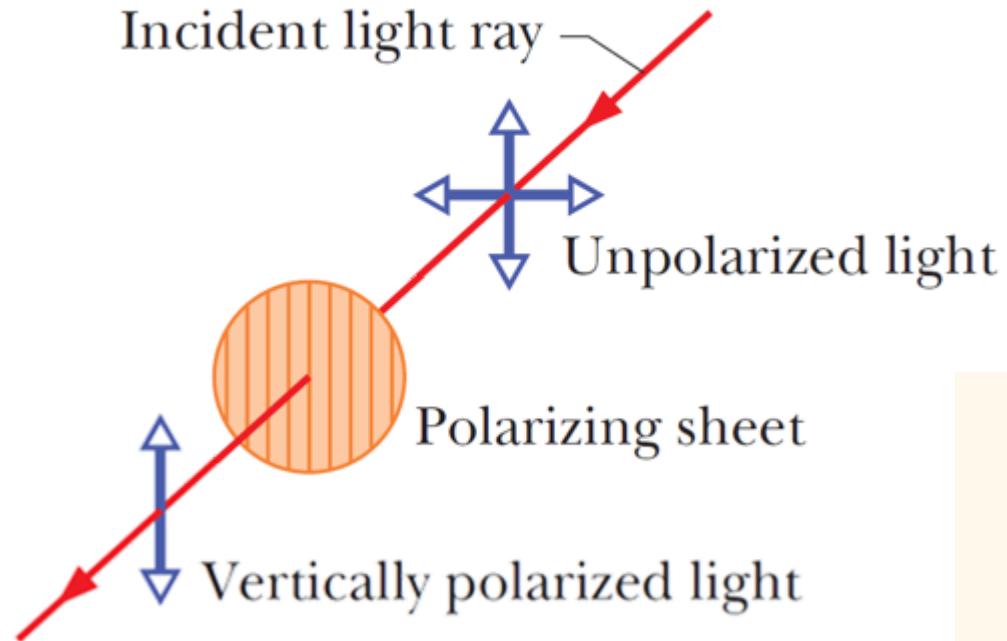


**Unpolarized → Polarized light**

Use of **polarizing sheet**

→ material that **transmit light**  
polarized along **one direction**  
**= Polarizing direction**

# POLARIZATION



Emerging light is linearly polarized parallel to the polarizing direction

**Unpolarized → Polarized light**

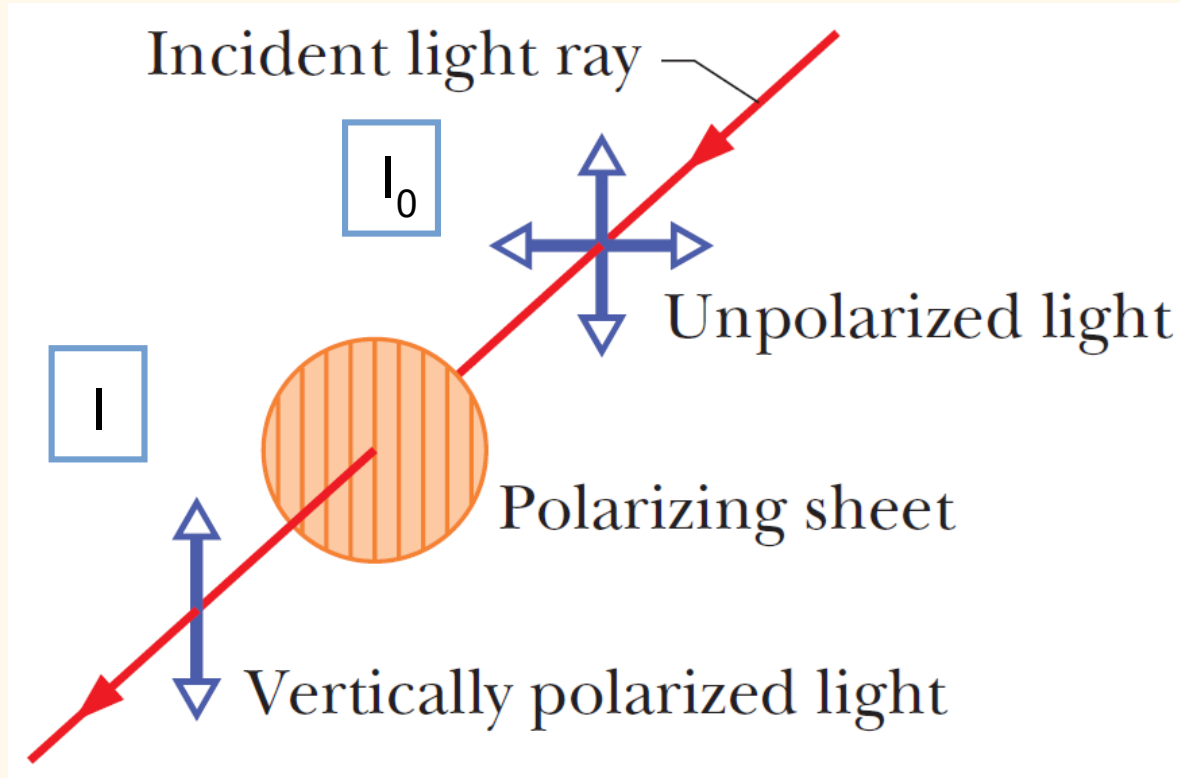
Use of **polarizing sheet**

→ material that **transmit light** polarized along **one direction**  
**= Polarizing direction**

→ **absorbs light** polarized along the **perpendicular direction**

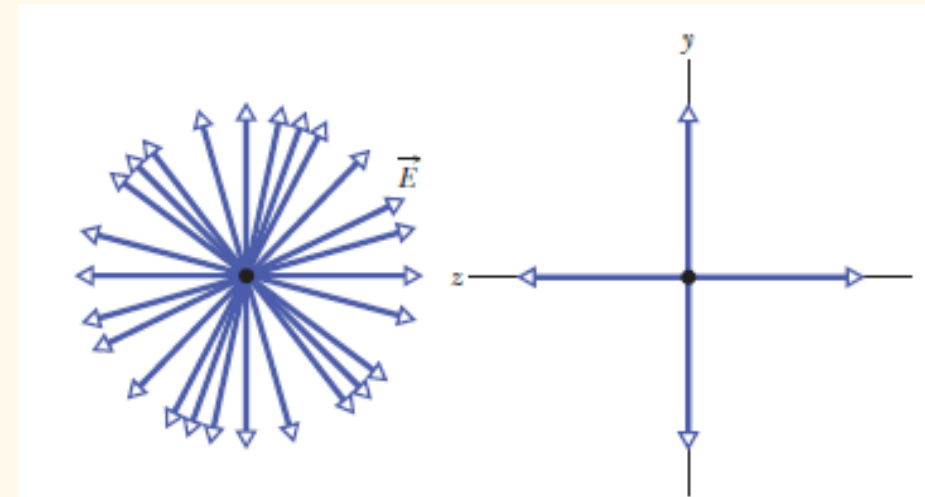
Note: like always in physics, it's more complicated than that. Polarizers don't really act like filters by « absorbing » some light. We have to deal with quantum mechanics for a full understanding

# POLARIZATION



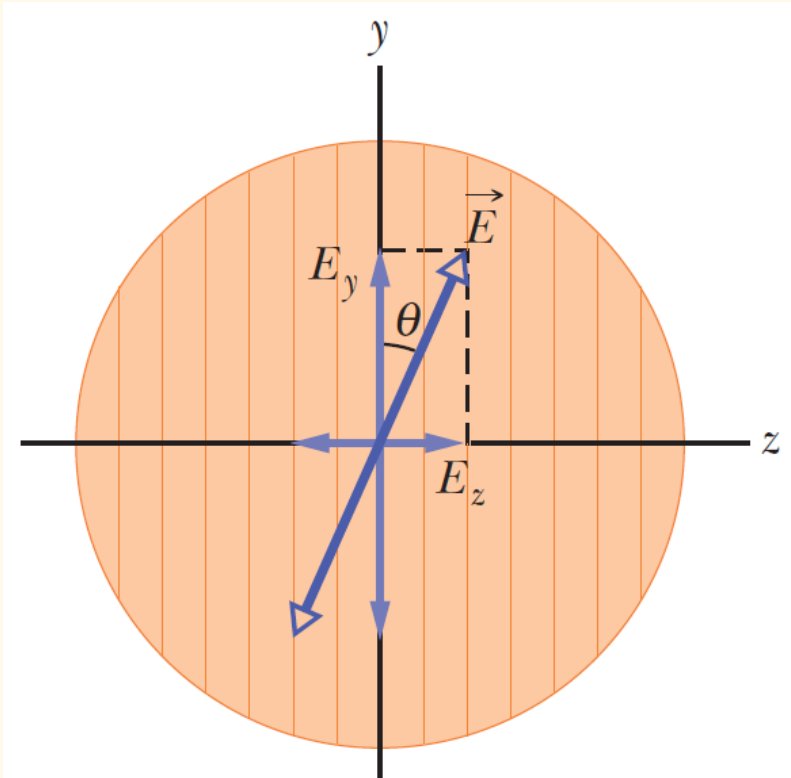
Intensity of the emerging light:

- For unpolarized light



$$I = I_0 / 2 \quad (\text{one-half rule})$$

# POLARIZATION



Note: 2 polarizers with perpendicular polarizing directions transmit no light

Intensity of the emerging light:

→ For linearly **incident** polarized light

$$\vec{E} = E_y \vec{u}_y + E_z \vec{u}_z = E \cos(\theta) \vec{u}_y + E \sin(\theta) \vec{u}_z$$

$$I_0 = \frac{E^2}{2c\mu_0}$$

→ For **Emerging** light (sheet polarized along Oy)

$$\vec{E} = E_y \vec{u}_y = E \cos(\theta) \vec{u}_y$$

$$I = \frac{E^2}{2c\mu_0} \cos^2(\theta)$$

$$I = I_0 \cos^2(\theta)$$

(Malus law)

# REFLECTION AND REFRACTION

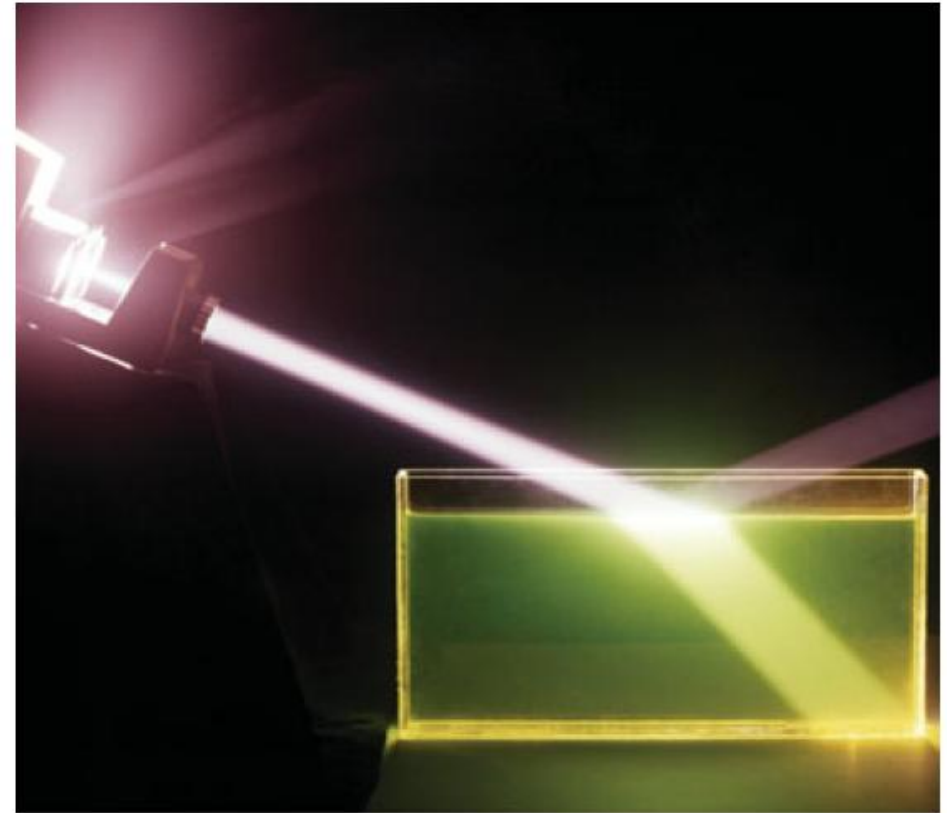
## Geometrical optics

We only consider **rays of light** and do not treat light as a wave (but rays have a wavelength)

→ **approximation**

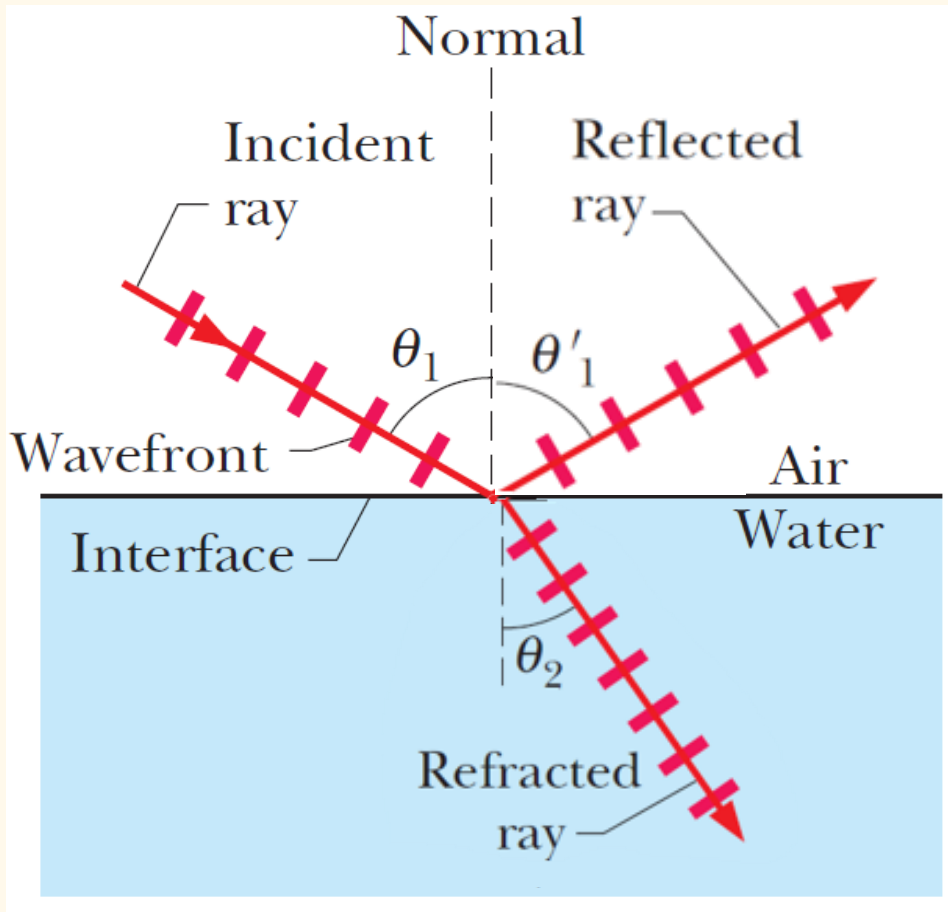
Easier to treat propagation of light:

- Interfaces between mediums
- Design of optical instruments



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# REFLECTION AND REFRACTION

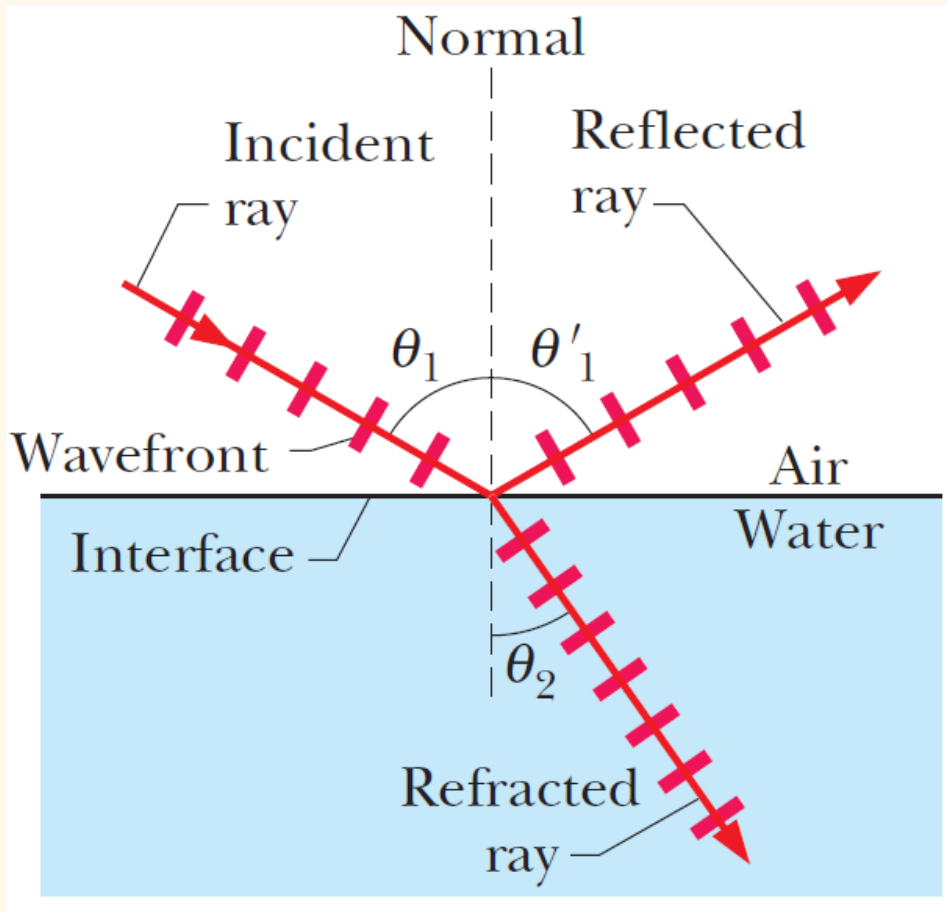


An **Incident ray** of light arrives at an interface (here air / water) With an angle  $\theta$  with respect to the **normal to the interface**

- Some light is reflected  
→ **Reflected ray**

- Some light travels in the other medium  
→ **Refracted ray**

# REFLECTION AND REFRACTION



## Vocabulary

$\theta_1$  is the **angle of incidence**

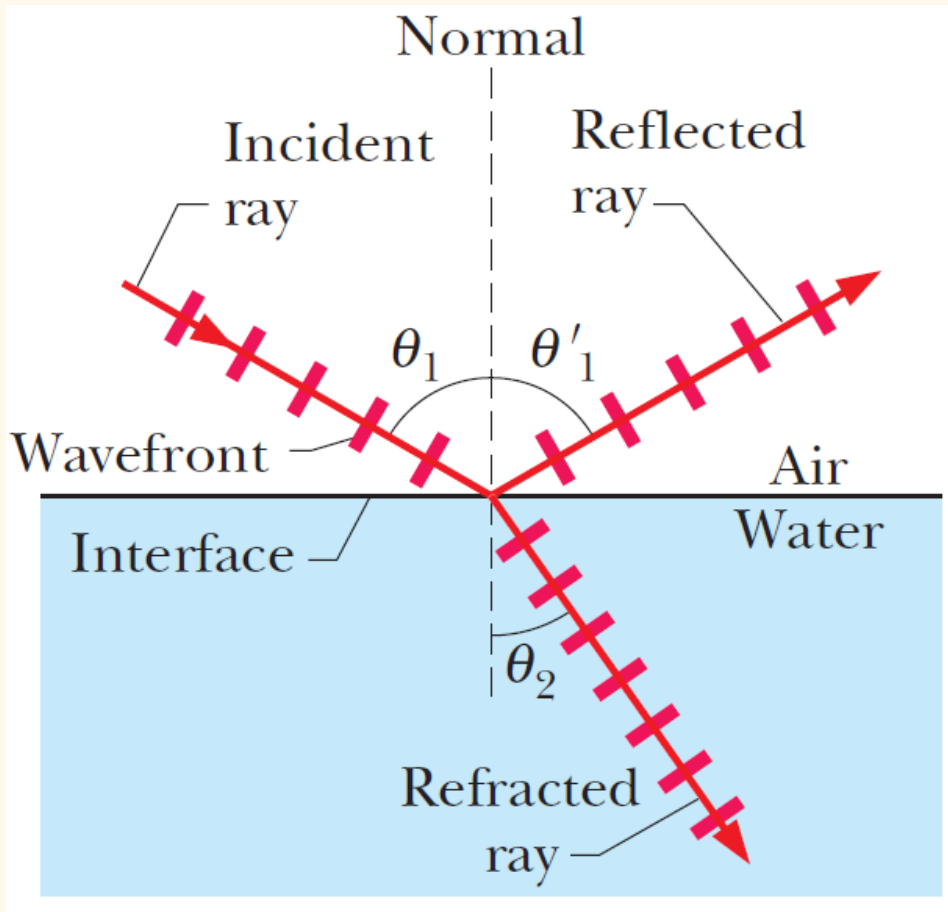
$\theta'_1$  is the **angle of reflection**

$\theta_2$  is the **angle of refraction**

**Angles between the rays and the normal**

The normal and the incident ray define the **plane of incidence**

# REFLECTION AND REFRACTION



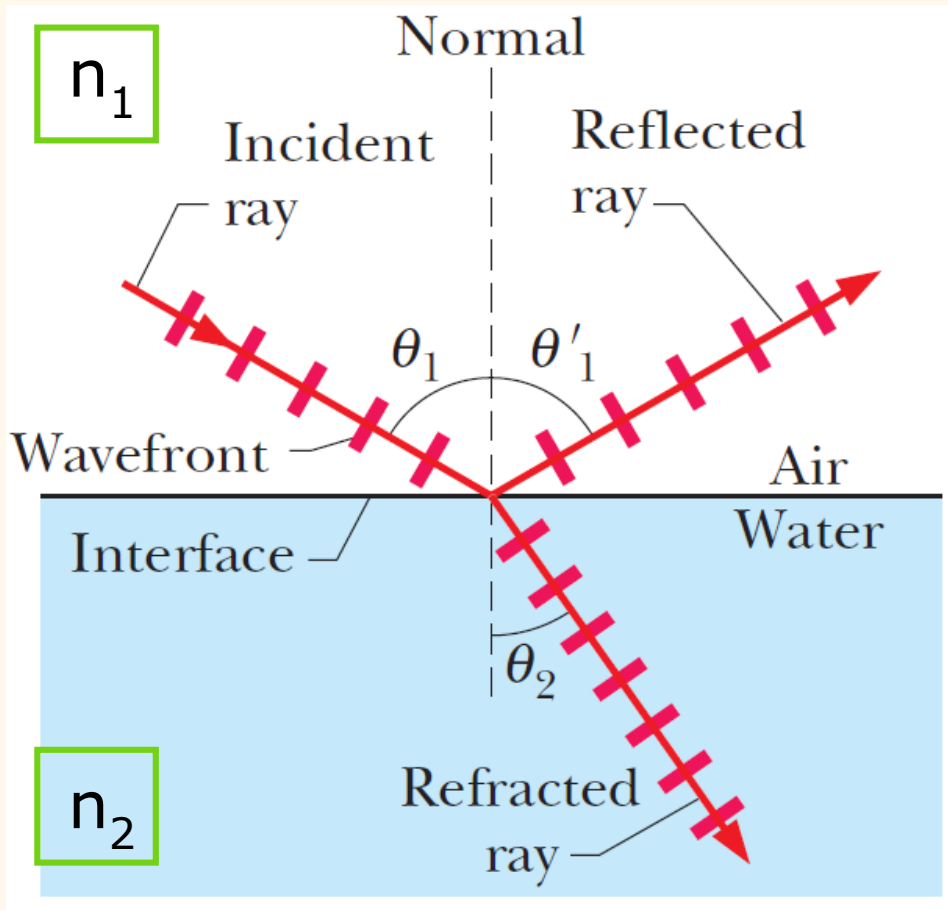
An **Incident ray** of light arrives at an interface (here air / water) With an angle  $\theta$  with respect to the **normal to the interface**

Directions of the reflected ray and the refracted ray are determined by Snell's laws

Note: In France we call them *The Laws of Snell – Descartes* (S-D)



# REFLECTION AND REFRACTION



## 1<sup>st</sup> S-D law:

The **reflected ray** and the **refracted ray** are in the **plane of incidence**

## 2<sup>nd</sup> S-D law:

The **angle of reflection** is equal to the **angle of incidence**

$$\theta'_1 = \theta_1$$

## 3<sup>rd</sup> S-D law:

The **angle of refraction** is given by

$$n_1 \sin(\theta_1) = n_2 \sin(\theta_2)$$

$n_1, n_2$ : **indexes of refraction** of the mediums

# REFLECTION AND REFRACTION

$$n = \frac{\text{speed of light in the vacuum}}{\text{speed of light in the medium}} = \frac{c}{v}$$

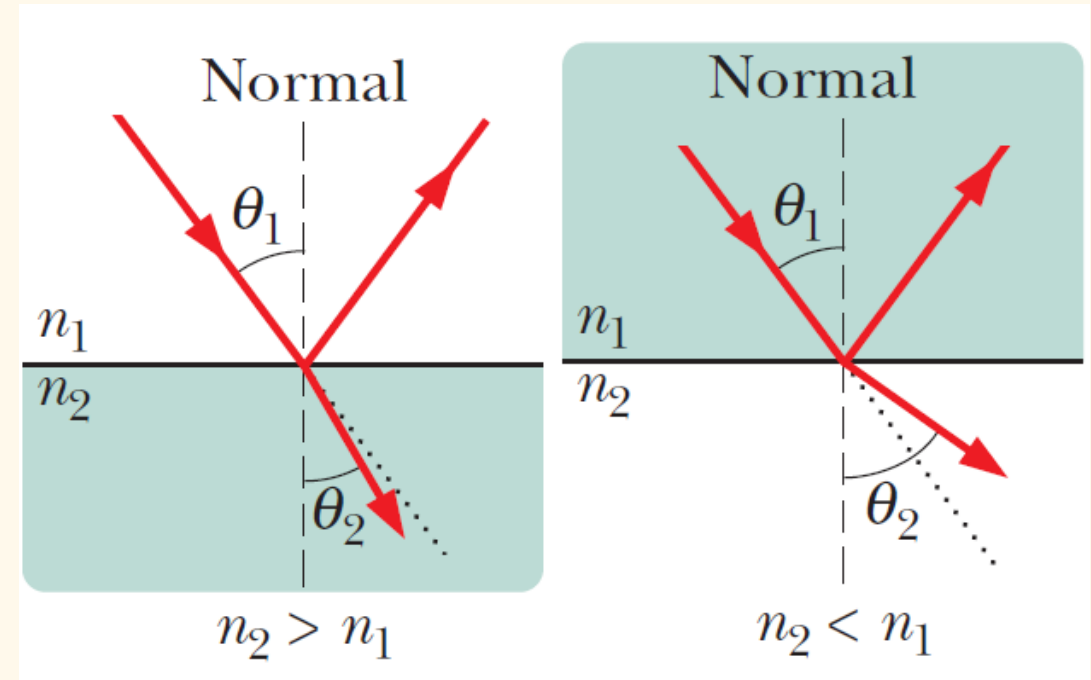
Typical values:  $\begin{cases} n_{\text{air}} \simeq 1 \\ n_{\text{eau}} \simeq 1.33 \\ n_{\text{glass}} \simeq 1.4 - 1.7 \end{cases}$

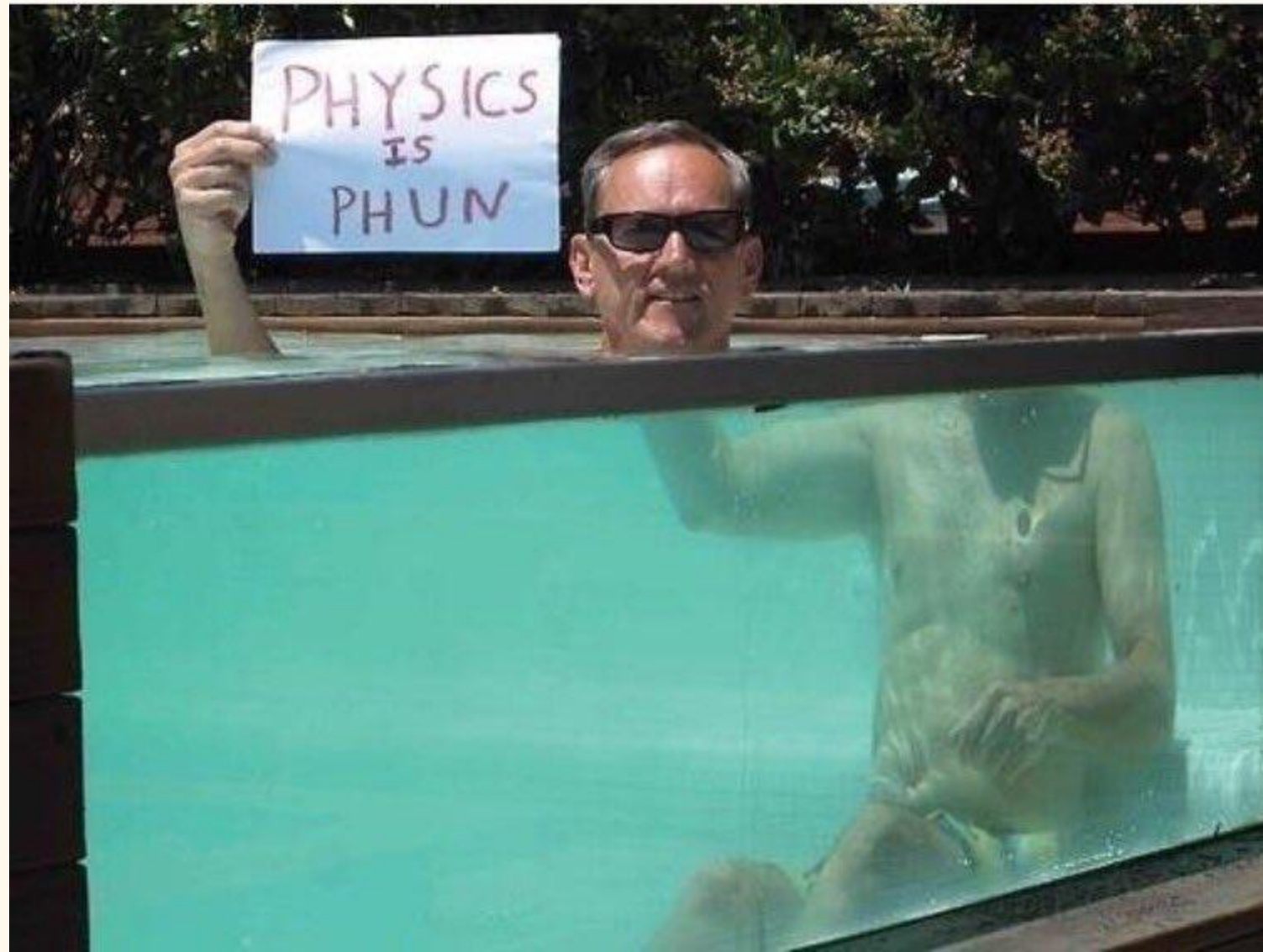
Rewriting the 3<sup>rd</sup> S-D law

$$\sin(\theta_2) = \frac{n_1}{n_2} \sin(\theta_1)$$

If  $n_2 > n_1$  (resp.  $< n_1$ )  
the refracted ray is bend towards  
(resp. away) from the normal

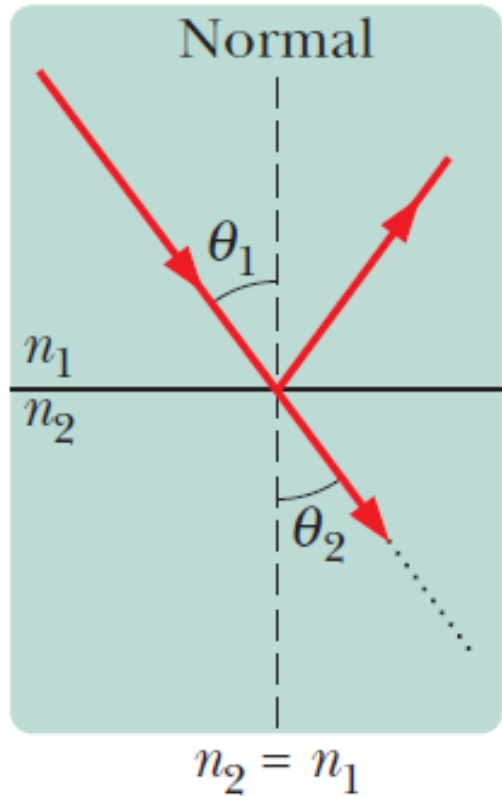
**That implies:  $n \geq 1$**



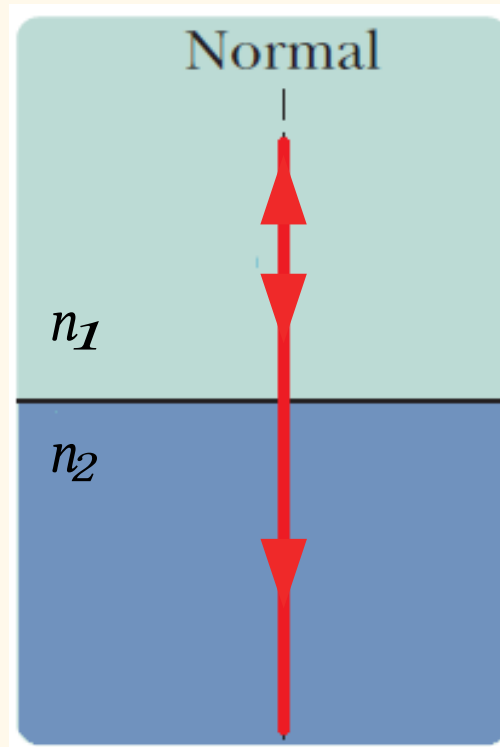


# REFLECTION AND REFRACTION

Note:



If  $n_1 = n_2$  (index matching) the incident ray is not deviated



If  $\theta_1 = 0^\circ$  (normal incidence)

$$\rightarrow \sin(\theta_1) = 0$$

Thus,  $\theta'_1 = 0$

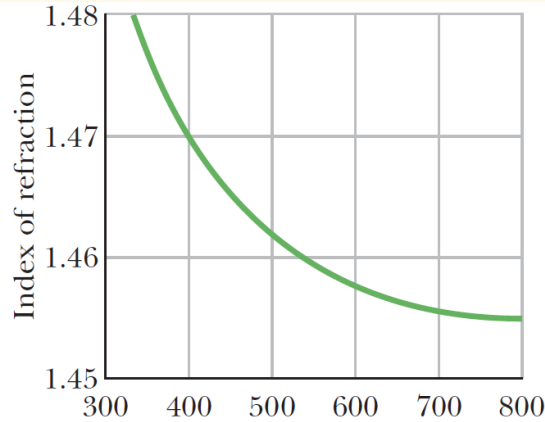
$\rightarrow$  The reflected beam is superposed to the incident ray

And  $\sin(\theta_2) = 0$ , so  $\theta_2 = 0$

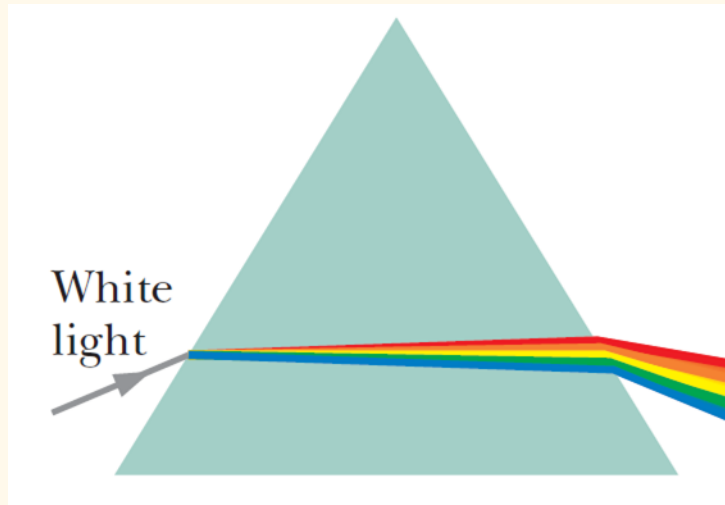
$\rightarrow$  The refracted ray is not deviated

# REFLECTION AND REFRACTION

Note that **n is a function of the wavelength** (e.g.  $n_{(\lambda)}$  for fused quartz)



Beam of white light = superposition of **Monochromatic** rays (each with a different  $\lambda$ ) with the same  $\theta_1$   
→ different  $\theta_2$   
→ **Chromatic dispersion**



Prisms separate wavelengths with two refractions

The same phenomena occurs in raindrops to form rainbows

# TOTAL INTERNAL REFLECTION

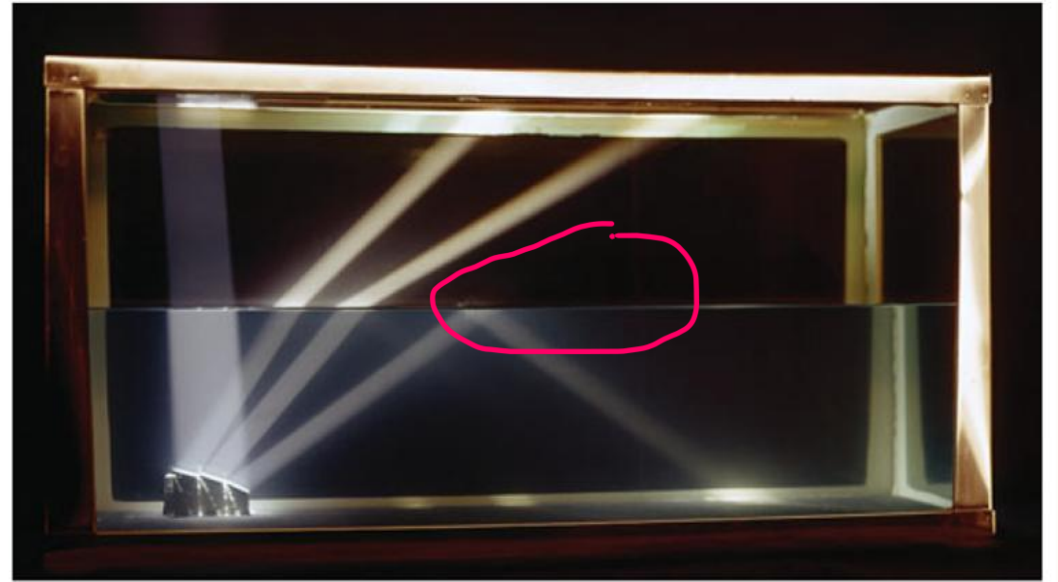
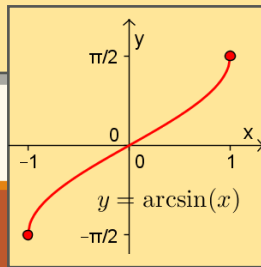
## S-D law of refraction:

$$n_1 \sin(\theta_1) = n_2 \sin(\theta_2)$$

$$\rightarrow \theta_2 = \arcsin\left(\frac{n_1}{n_2} \sin(\theta_1)\right)$$

We assume  $n_1 > n_2$   
(e.g. medium 1 is water, 2 is air)

→ Above a certain value of  $\theta_1$ , we cannot calculate  $\theta_2$   
( $\arcsin(x)$  with  $x > 1$  is not defined)



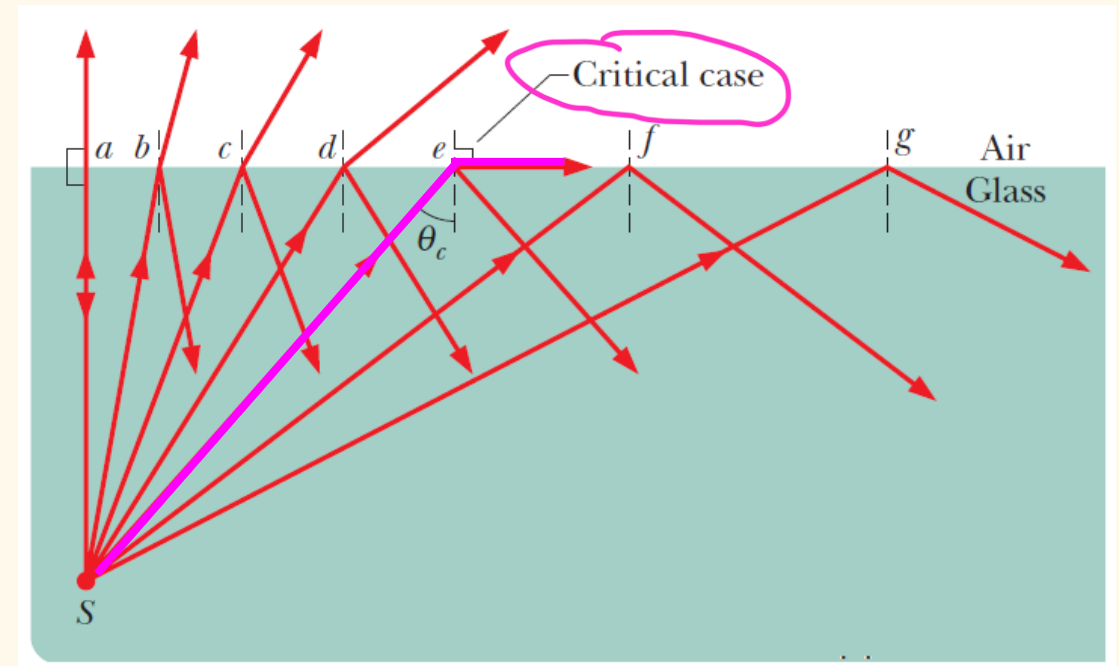
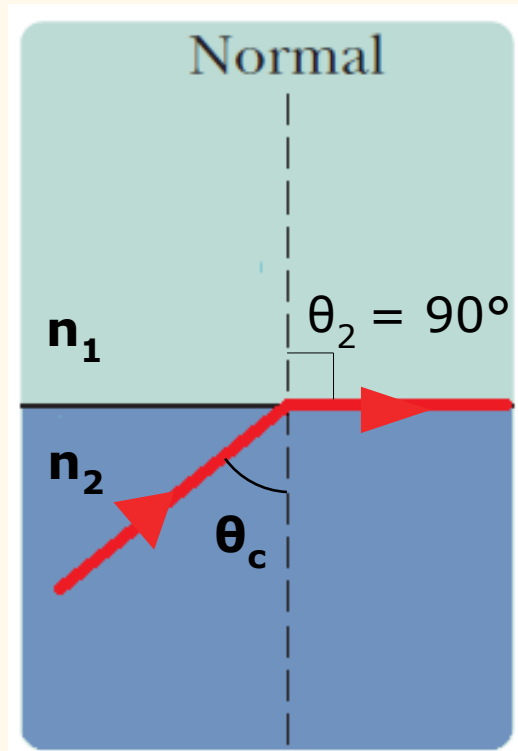
Ken Kay/Fundamental Photographs

In such cases, the refracted ray **does not exist** → only reflection in medium 1  
→ **Total Internal Reflection (TIR)**

# TOTAL INTERNAL REFLECTION

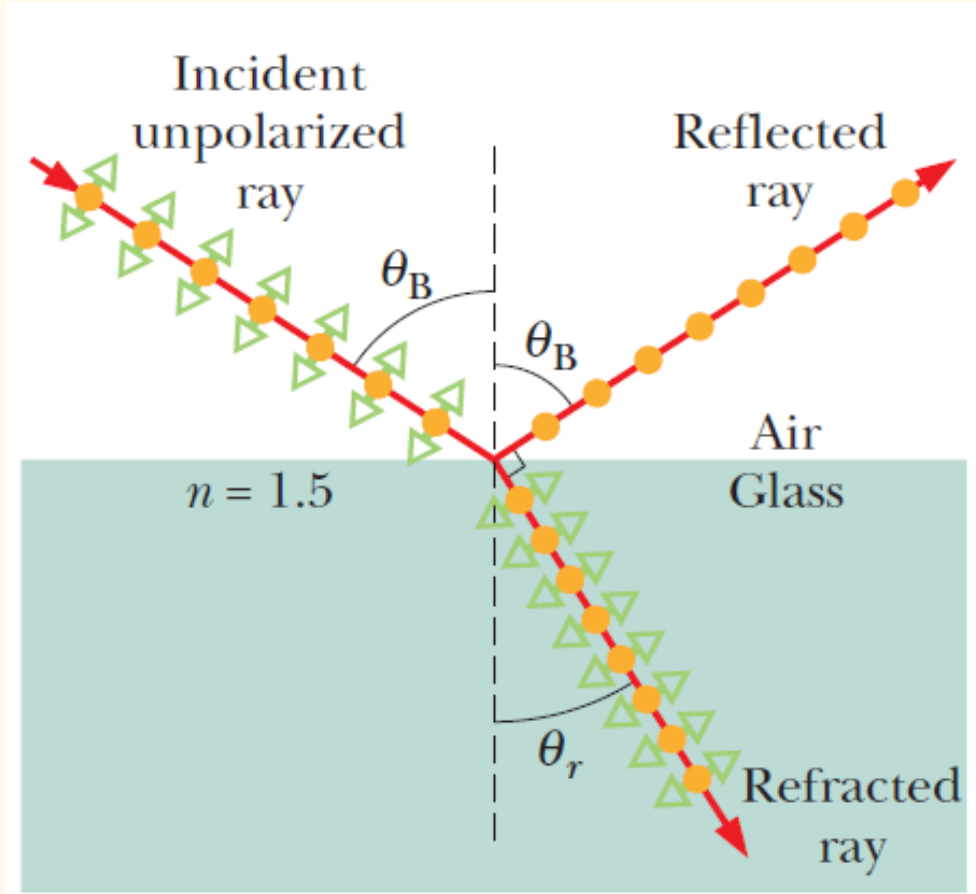
For TIR, we define the **critical angle**  $\theta_c$  such that  $\theta_2 = 90^\circ$ :  
 $n_1 \sin(\theta_c) = n_2 \sin(90) = n_2$

$$\theta_c = \text{asin}\left(\frac{n_2}{n_1}\right)$$



**TIR if  $n_1 > n_2$  and  $\theta > \theta_c$**

# POLARIZATION BY REFLECTION



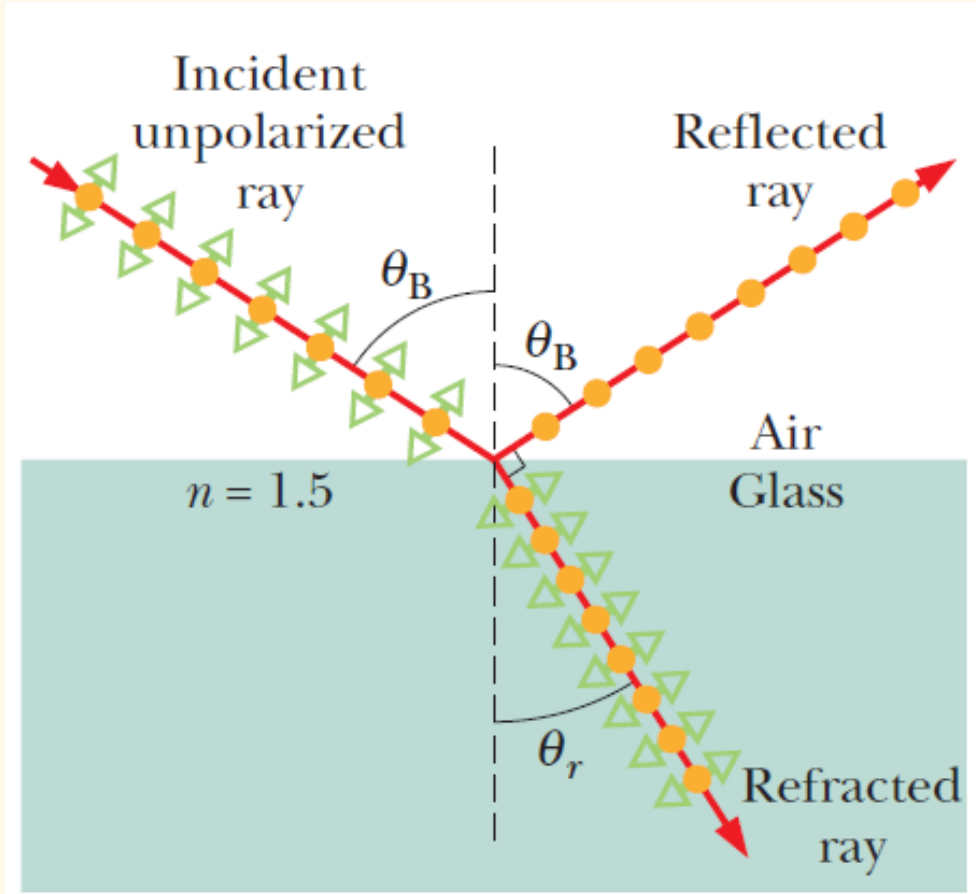
Previously, we saw that light passing through a polarizing sheet is polarized

**Reflection Brewster angle  $\theta_B$**   
For  $\theta_1 = \theta_B$  (2<sup>nd</sup> S-D)

- **Reflected and refracted rays are perpendicular**
- **Reflected ray is polarized perpendicularly to the plane of incidence**



# POLARIZATION BY REFLECTION



We have:  $\theta_B + \theta_2 = 90$

with  $n_1 \sin(\theta_B) = n_2 \sin(\theta_2)$  (3<sup>rd</sup> S-D)

So:  $n_1 \sin(\theta_B) = n_2 \sin(90 - \theta_B) = n_2 \cos(\theta_B)$

Then:  $\tan(\theta_B) = \frac{n_2}{n_1}$  ———  $\theta_B = \arctan\left(\frac{n_2}{n_1}\right)$

One of my research interest was to retrieve the Snell laws and total Internal Reflexion, not with light but with high speed electrons, called ballistic electrons that « behave like light » in electronics 2D based material

laboratoire pierre algrain  
électronique et photonique quantiques

Les Rencontres de Moriand, March 2019

GRAPHENE FLAGSHIP

**A corner reflector for graphene Dirac fermions as a phonon-scattering sensor**

arXiv:1901.02225 (2019)



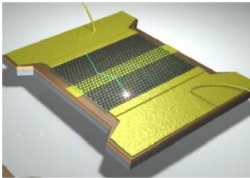
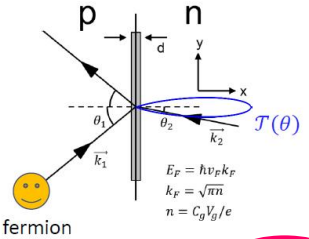
David Mele

H. Graef<sup>1,2</sup>, Q. Wilmart<sup>1</sup>, M. Rosticher<sup>1</sup>, D. Mele<sup>1</sup>, L. Banszerus<sup>1</sup>, C. Stampfer<sup>1</sup>,  
T. Taniguchi<sup>3</sup>, K. Watanabe<sup>3</sup>, E. Bocquillon<sup>1</sup>, G. Féa<sup>1</sup>, J.M. Berroir<sup>1</sup>, E.H.T. Teo<sup>1</sup> and S. Plaçais<sup>1</sup>

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ENS PSL CIMS UPMC NANYANG UNIVERSITY RWTH AACHEN UNIVERSITY NIMS CNRS

Dirac Fermion Optics (DFO) – Diopter

fermion

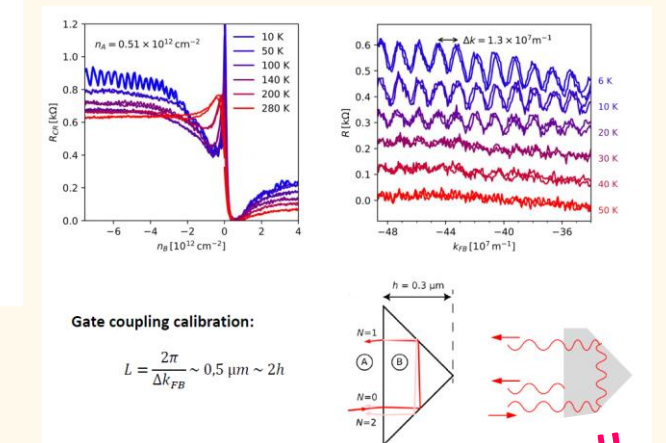
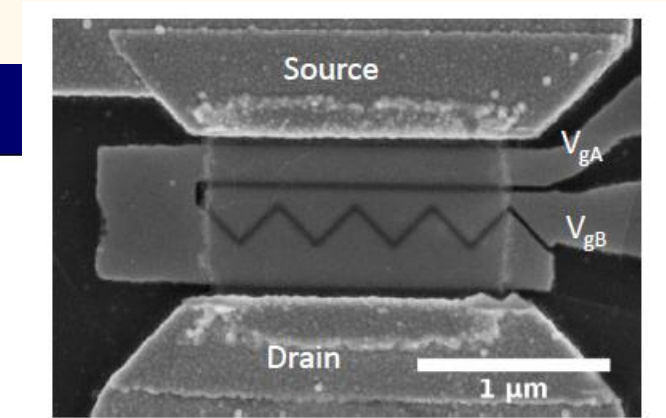
	Photon optics (3D)	Dirac fermion optics (2D)
Medium	transparent	ballistic
Phase velocity	$3 \cdot 10^8 \text{ m s}^{-1}$	$10^6 \text{ m s}^{-1}$
Snell-Descartes	$n_1 \sin \theta_1 = n_2 \sin \theta_2$	$E_{F1} \sin \theta_1 = E_{F2} \sin \theta_2$
Critical angle	$\theta_c = \arcsin\left(\frac{n_2}{n_1}\right)$	$\theta_c = \arcsin\left(\frac{E_{F2}}{E_{F1}}\right)$
Fresnel relation	$R_s = \left  \frac{n_1 \cos \theta_i - n_2 \cos \theta_t}{n_1 \cos \theta_i + n_2 \cos \theta_t} \right ^2$	$T(\theta) = e^{-\pi \frac{2d}{ k_1 - k_2 } k_2^2 \sin^2 \theta}$

Chelarov et al., PRB 74 (2006) 041403(R)

J. Cayssol et al., PRB 79 (2009) 075428

Light in two different  
optical medium

electron in two different  
electrostatic medium

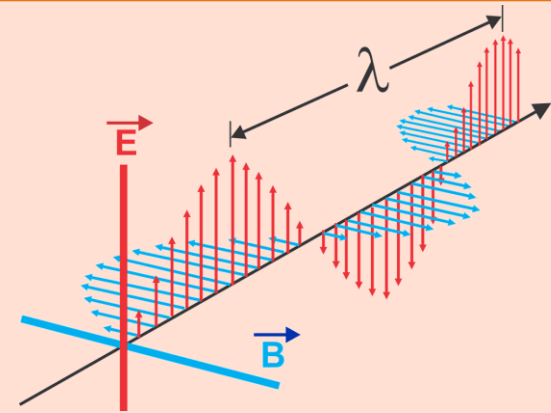


TIR with electron confirmed!

# KEY POINTS

Light as an EM wave

Structure of the traveling plane EM wave



Direction of propagation

$$\frac{E}{B} = c$$

Poynting vector, intensity and radiation pressure

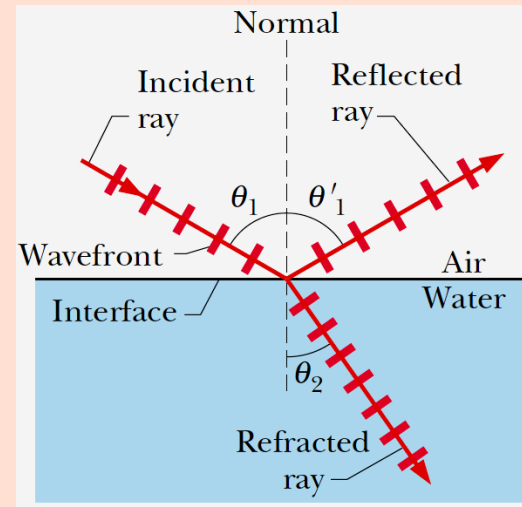
$$\vec{S} = \frac{\vec{E} \times \vec{B}}{\mu_0}$$

$$I = \frac{E_m^2}{2c\mu_0}$$

Polarized and unpolarized light

Snell-Descartes laws

TIR and Polarization by reflection



$$n_1 \sin(\theta_1) = n_2 \sin(\theta_2)$$

# READING ASSIGNMENT

**Chapter 34 of the textbook**